Quantum amplitudes for circular polarization

The field of any plane light wave with frequency v can be written as $\vec{E} = \operatorname{Re}\left[\left(\frac{\vec{e}_x - i\vec{e}_y}{\sqrt{2}}B_R + \frac{\vec{e}_x + i\vec{e}_y}{\sqrt{2}}B_L\right)e^{i(kz-\omega t)}\right]$

The Intensity in the **R** channel after an **R/L** analyzer:

$$I_R \propto 2 \left\langle \operatorname{Re}\left[\frac{\vec{e}_x - i\vec{e}_y}{\sqrt{2}} B_R e^{i(kz - \omega t)}\right]^2 \right\rangle_t$$

$$= 2 \left\langle \frac{1}{4} \left[2B_R B_R^* \frac{\vec{e}_x - i\vec{e}_y}{\sqrt{2}} \frac{\vec{e}_x + i\vec{e}_y}{\sqrt{2}} \right] \right\rangle_t = B_R B_R^*$$

Intensity in the L channel: $I_L \propto B_L B_L^*$

The corresponding state vector describes a photon which is in the state $|R\rangle$ with probability $|B_R|^2$ and in the state $|L\rangle$ with probability $|B_L|^2$

The interferences of the photons correspond to the interferences of the field when the phases and the probability amplitudes are chosen according to the amplitudes and phases of the filed components by using complex quantum amplitudes

 $|\Psi >= |R > B_R + |L > B_L$

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03/18/2005



Formal correspondence to polarization states

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03/18/2005

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Polarization statesStationary states of the wave function
Orthogonality $\langle x | y \rangle = 0$, $\langle R | L \rangle = 0$ $\int_{0}^{\infty} \Phi_n^*(x) \Phi_m(x) dx = \delta_{nm}$

Completeness: Every states can be described in terms of the basis states.

 $|\Psi\rangle = |x\rangle A_x + |y\rangle A_y$ $f(x) = \sum_{n=0}^{N} \Phi_n(x) A_n$

The expansion amplitudes are found by projections; they are projection amplitudes.

$$A_{x} = \left\langle x \middle| \Psi \right\rangle \qquad \qquad A_{n} = \int_{-\infty}^{\infty} \Phi_{n}^{*}(x) f(x) \, dx$$

Due to this correspondence, also the stationary states of the wave function are written as **ket vector**:

$$\Phi_{n}(x) \rightarrow |n\rangle \text{ with } \langle n|m\rangle = 0$$

$$|f\rangle = \sum_{n=0}^{N} |n\rangle A_{n} \text{ with } A_{n} = \langle n|f\rangle$$

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