

P331B Section

29 March 2013

TODAY: CANONICAL TRANSFORMATIONS

RETURN HW

NEXT WK: MIDTERM REVIEW

↳ bring your questions

RECALL (HW #2, i think):

↳ LAGRANGIAN

LEOM ARE UNCHANGED UNDER THE TRANSFORMATION

$$\boxed{L' = \lambda L - \frac{dF}{dt}} \quad (*)$$

↑
RESCALING
(CANCELS ON
BOTH SIDES
of EOM)
ignore this.

↑
total time derivative

BECOMES BOUNDARY TERM
IN ACTION, BUT IS DON'T
VARY ON THE BOUNDARY.

recall further: $\boxed{H = p\dot{q} - L}$ (implied ϵ)

~~GOAL~~

GOAL: WANT TO FIND NEW COORDINATES (P, Q) IN
THE HAMILTONIAN PICTURE SUCH THAT
HEOM ARE UNCHANGED

↳ "canonical transformation"

Func of p, q

↳

RHS

EXPLICITLY: WANT $Q(q, p) + P(q, p)$ s.t. $\exists K$,
A NEW HAMILTONIAN, SUCH THAT:

$$\dot{Q}_i = \frac{\partial K}{\partial P_i} \quad \dot{P}_i = - \frac{\partial K}{\partial Q_i}$$

IF THESE ARE GOOD COORDINATES, THEN THE
LAGRANGIAN ASSOCIATED TO K MUST BE AT MOST
DIFFERENT FROM L (ORIGINAL LAGRANGIAN) BY A
TRANSFORMATION OF THE FORM (*)

$$L' = \boxed{P\dot{Q} - K = \lambda (P\dot{Q} - H) - \frac{dF}{dt}}$$

↑
L

for our purposes, set $\lambda = 1$

USUALLY VERY
INTERESTING

(scale transformations are uninteresting
in this context.)

← RHS A PRIORI
UNKNOWN

so: $\boxed{P\dot{Q} - H = P\dot{Q} - K + \frac{dF}{dt}}$

Philosophy:
DON'T KNOW
P, Q, K ...
JUST SPECIFY F

F: GENERATING FUNCTION, CHARACTERIZES THE TRANSF.
→ P, Q, K

simplest example: $F = F_1(q, Q, t)$

OTHER CASES: q, P
 p, Q
 p, P

↑ require additional
terms

$$\begin{aligned}
 \text{then: } p\dot{q} - H &= P\dot{Q} - K + \frac{\partial F_1}{\partial t} \\
 &= P\dot{Q} - K + \underbrace{\frac{\partial F_1}{\partial t} + \frac{\partial F_1}{\partial q}\dot{q} + \frac{\partial F_1}{\partial Q}\dot{Q}}_{\text{came from } F_1(q, Q, t)}
 \end{aligned}$$

\uparrow $H(p, q)$ \uparrow $K(P, Q)$

note: q & Q are separately independent
 (@ this point we haven't specified $Q(p, q)$)

So: coefficient of \dot{q} on LHS & RHS MATCH:

$$\boxed{P = \frac{\partial F_1}{\partial q}} \rightarrow \text{CAN INVERT TO SOLVE FOR } Q(q, p) \text{ hopefully!}$$

ALSO: coefficient of \dot{Q} on LHS & RHS MATCH:

$$0 = P + \frac{\partial F_1}{\partial Q} \Rightarrow \boxed{P = -\frac{\partial F_1}{\partial Q}}$$

FURTHER: WHAT'S LEFT OVER MUST BE EQUAL:

$$-H = -K + \frac{\partial F_1}{\partial t} \Rightarrow \boxed{K = H + \frac{\partial F_1}{\partial t}}$$

SO: GIVEN $F_1(q, Q)$, WE HAVE FOUND THE
 P , Q (implicitly), and K !

SIMPLE EXAMPLE

$$F_1 = gQ \Rightarrow \begin{cases} p = Q \\ P = -g \\ K = H \end{cases} \quad \leftarrow \text{trivially gives } Q$$

$$\Rightarrow \begin{cases} Q = p \\ P = -g \end{cases} \quad \left. \vphantom{\begin{cases} Q = p \\ P = -g \end{cases}} \right\} \text{swapped } p \text{ ' } g!$$

↑
PHYSICALLY THIS IS WEIRD!

~ ROTATING PHASE SPACE

$$H(p, g) \mapsto H(-g, p)$$

OTHER TYPES OF CANONICAL TRANSFORMATIONS
→ by LEGENDRE TRANSFORMS

$$\text{eg. } F_3(p, Q, t) = F_1(q, Q, t) - \cancel{g} p$$

↑
LEGENDRE TRANSFORM $q \rightarrow p$
note: not the usual $q \rightarrow p$!

$$F = F_1(q, Q, t) = \boxed{F_3(p, Q, t) + gP}$$

$$\text{eg. } F_2(q, P, t) = F_1(q, Q, t) + QP$$

$$F = F_1(q, Q, t) = \boxed{F_2(q, P, t) - QP}$$

note sign!

SUMMARY :

$$F = F_1(g, Q, t)$$

$$F = F_2(g, P, t) - QP$$

$$F = F_3(p, Q, t) + gP$$

$$F = F_4(p, P, t) + gP - QP$$

Why the signs? try wrong sign:

$$F = F_2(g, P) \pm QP$$

$$\Rightarrow \underbrace{p\dot{g}} - H = \underbrace{P\dot{Q}} - K + \frac{\partial F_2}{\partial t} + \underbrace{\frac{\partial F_2}{\partial g}} \dot{g} + \left(\frac{\partial F_2}{\partial P} \pm Q \right) \dot{P} \pm \underbrace{P\dot{Q}}$$

$$\text{left w/ : } -H = -K + \frac{\partial F_2}{\partial t} + P\dot{Q} \pm P\dot{Q}$$

GOOD, WE
BEFORE

WANT THIS TO VANISH

DONT WANT K

TO BE A FUNCTION

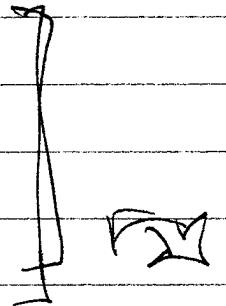
OF Q !

$$P = \frac{\partial F_2}{\partial g}$$

$$Q = \frac{\partial F_2}{\partial P}$$

~~K = H + \dots~~

$$K = H + \frac{\partial F_2}{\partial t}$$



Simple example

$$F_2 = \mathcal{F}P \quad \rightarrow \quad P = \frac{\partial F_2}{\partial \mathcal{F}} = \mathcal{F} \quad \left. \vphantom{\frac{\partial F_2}{\partial \mathcal{F}}} \right\} \text{trivial transf.}$$
$$Q = \frac{\partial F_2}{\partial P} = \mathcal{F}$$
$$K = H$$

Non-trivial example

$$\omega^2 = k/m$$

$$H = \frac{1}{2m} (p^2 + m^2 \omega^2 q^2)$$

for simplicity: ~~$m = \omega = 1$~~ (dimensions...)
 $m = 1/2 \quad \omega = 2$

$$H = \underbrace{p^2 + q^2}$$

looks LIKE A CIRCLE!

WANT:

$$p = f(P) \cos Q$$
$$q = f(P) \sin Q$$

$$K = f^2(P) (\cos^2 Q + \sin^2 Q) = f^2(P)$$

aha! Q is a cyclic (ignorable) coordinate!

How to get ψ 's?

eg: $F_1 = g^2 \cos Q$

$$P = \frac{\partial F_1}{\partial \dot{Q}} = 2g \cos Q$$

$$\psi = -\frac{\partial F_1}{\partial Q} = g^2 \frac{1}{\sin^2 Q}$$

$$g = \sqrt{\psi} \sin Q$$
$$P = \sqrt{\psi} \cos Q$$

$$H = P^2 + \psi^2 = \psi$$

restoring $m \dot{\omega}$: $P = E/\omega$

$$\dot{Q} = \omega \rightarrow Q = \omega t + Q_0$$

$$g = \sqrt{\frac{2E}{m\omega^2}} \sin(\omega t + Q_0)$$

$$P = \sqrt{2mE} \cos(\omega t + Q_0)$$

→ REVIEW!

An example of Poisson Brackets:

$$\{f, g\} = \sum_i \frac{\partial f}{\partial q_i} \frac{\partial g}{\partial p_i} - \frac{\partial f}{\partial p_i} \frac{\partial g}{\partial q_i}$$

Why is it useful?

$$\dot{A} = \frac{\partial A}{\partial t} + \frac{\partial A}{\partial q} \dot{q} + \frac{\partial A}{\partial p} \dot{p}$$

$\uparrow \qquad \qquad \uparrow$
 $\frac{\partial H}{\partial p} \qquad - \frac{\partial H}{\partial q}$

$$= \frac{\partial A}{\partial t} + \{A, H\}$$

LOOKS A LOT LIKE
QM! H "GENERATES"
TIME TRANSLATIONS.

CONSIDER: $V(r) = -\frac{k}{r}$

KEPLER PROBLEM

CLAIM: RUNGE LENZ VECTOR IS CONSERVED

$$\vec{A} \equiv \vec{p} \times \frac{\vec{L}}{m} - \frac{k}{r} \hat{r}$$

① WHAT IS THE HAMILTONIAN?

$$L = \frac{1}{2} m(\dot{r})^2 + \frac{1}{2} m(r^2 \dot{\theta})^2 - \frac{k}{r}$$

$$p_r = m\dot{r} \qquad l_z = m r^2 \dot{\theta}$$

$$H = \frac{p_r^2}{2m} + \frac{l_z^2}{2mr^2} - \frac{k}{r}$$

Poisson Bracket facts

$$\bullet \quad \{ab, c\} = a\{b, c\} + \{a, c\}b$$

$$\begin{aligned} \frac{\partial(ab)}{\partial q} \frac{\partial c}{\partial p} - (q \leftrightarrow p) &= \left(a \frac{\partial b}{\partial q} \frac{\partial c}{\partial p} + b \frac{\partial a}{\partial q} \frac{\partial c}{\partial p} \right) - (q \leftrightarrow p) \\ &= a \left(\frac{\partial b}{\partial q} \frac{\partial c}{\partial p} - (q \leftrightarrow p) \right) + \left(\frac{\partial a}{\partial q} \frac{\partial c}{\partial p} - (q \leftrightarrow p) \right) b \end{aligned}$$

$$\bullet \quad \vec{l} = \vec{r} \times \vec{p} \rightarrow l_i = \epsilon_{ijk} r_j p_k$$

$$\begin{aligned} \{l_i, r_a\} &= \epsilon_{ijk} \{r_j p_k, r_a\} \\ &= \epsilon_{ijk} (r_j \{p_k, r_a\} + \underbrace{\{r_j, r_a\}}_{=0} p_k) \\ &= \epsilon_{ijk} r_j (-\delta_{ka}) \\ &= \epsilon_{iaj} r_j \end{aligned}$$

$$\text{similarity } \{l_i, p_a\} = \pm \epsilon_{ike} p_e$$

MORE USEFUL

$$\begin{aligned}\{ \vec{l}_i, r^2 \} &= \{ l_i, \sum_a r_a^2 \} = \sum_a 2r_a \{ l_i, r_a \} \\ &= 2r_a (\epsilon_{ias} r_s) \\ &= 2\epsilon_{ias} r_a r_s = 0\end{aligned}$$

similarly $\{ l_i, p^2 \} = 0$

2. \vec{l} is conserved $H = H(p^2, r^2)$

$$\{ l_i, H \} = \sum_j \frac{\partial l_i}{\partial p_j} \frac{\partial H}{\partial r_j} - \left(\frac{\partial l_i}{\partial r_j} \frac{\partial H}{\partial p_j} \right) \quad (p \leftrightarrow r)$$

$$= \sum_j \frac{\partial l_i}{\partial p_j} \frac{\partial H}{\partial r^2} \frac{\partial r^2}{\partial r_j} - \left(\frac{\partial l_i}{\partial r_j} \frac{\partial H}{\partial p_j} \right)$$

but: $\frac{\partial l_i}{\partial p_j} \frac{\partial H}{\partial r^2} = \frac{\partial}{\partial p_j} \epsilon_{iab} r_a p_b$
 $= \epsilon_{iaj} r_a$

$$\text{so: } \frac{\partial l_i}{\partial p_j} \frac{\partial H}{\partial r^2} (2r_j) \sim \epsilon_{iaj} r_a r_j = 0$$

SIMILAR FOR $\frac{\partial l_i}{\partial r_j} \frac{\partial H}{\partial p^2} (2p_j)$

EXERCISE : FOR ~~THE~~ THE KEPLER PROBLEM,
SHOW THAT THE RUNGE-LENZ VECTOR
IS CONSERVED :

$$\vec{A} = \vec{p} \times \vec{l} + \frac{m\vec{r}k}{r}$$

$$\frac{\partial \Delta}{\partial t} = 0$$

~~$$\vec{A} = \vec{p} \times \vec{l} + \frac{m\vec{r}k}{r}$$~~

WE KNOW: $\{f(r), \vec{l}\} = \{p^2, \vec{l}\} = 0$

CAN PROVE $\frac{p^2}{2m} + \frac{k}{2mr^2} \rightarrow \frac{p^2}{2m}$

set $m=1$
 $k=1$

SINCE $p_z = 0$.

$$\dot{A} = \{H, A_k\} = \left\{ \frac{1}{2} p^2 + \frac{1}{r}, \epsilon_{kij} p_i l_j + \frac{r_k}{r} \right\}$$

$$= \frac{1}{2} \epsilon_{kij} \{p^2, p_i l_j\} + \epsilon_{ijk} \left\{ \frac{1}{r}, p_i l_j \right\} + \frac{1}{2} \{p^2, \frac{r_k}{r}\} + \left\{ \frac{1}{r}, \frac{r_k}{r} \right\}$$

$$\{p_j \{p^2, p_i\} + p_i \{p^2, l_j\}\} = 0$$

$$= \frac{1}{2} \{p^2, \frac{r_k}{r}\} + \epsilon_{ijk} \left\{ \frac{1}{r}, p_i l_j \right\}$$

$$\sum_a \frac{\partial(p^2)}{\partial r_a} \frac{\partial(r_k/r)}{\partial p_a} - \frac{\partial(p^2)}{\partial p_a} \frac{\partial(r_k/r)}{\partial r_a} = \sum_a -2p_a \cdot \left(\frac{\delta_{ka}}{r} - \frac{r_k r_a}{r^3} \right)$$

$$\frac{\partial a r_k}{r} + r_k \frac{\partial a}{\partial a} (r_k r_a)^{-1/2}$$

for example: let $k = x$ direction:

$$\left\{ P^2, \frac{\partial x}{\partial r} \right\} = -2P_x \left(\frac{1}{r} - \frac{x^2}{r^3} \right) + 2P_y \left(\frac{xy}{r^3} \right) + 2P_z \left(\frac{xz}{r^3} \right)$$

$$= \frac{-2P_x (y^2 + z^2)}{r^3} + \frac{2x (P_y y + P_z z)}{r^3}$$

claim: the $\frac{1}{r} P_i l_j$ term cancels this