Muon (g-2) Past and Future

With an emphasis on Beam Dynamics in the (g-2) Storage Ring

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The New g-2 Experiment:

A Proposal to Measure the Muon Anomalous Magnetic Moment to ± 0.14 ppm Precision

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http://lss.fnal.gov/archive/test-proposal/0000/fermilab-proposal-0989.shtml



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Fermilab P989

• First submitted to PAC in February 2009

- Great physics, how much does it cost?

- Building etc. costing exercise summer 09, report to PAC in November 2009
 - Great physics, Director must talk with DOE to figure out how to fund it.
- DOE Intensity Frontier Review August 2010
 - result embargoed until Presidential Budget released



Outline of the Talk

- Brief review of magnetic moments -- including the theory of a_{μ}
- Spin motion in a magnetic field
- Overview of the experimental technique
 - The precision storage ring magnet
 - The fast muon kicker
 - The electrostatic quadrupoles
- Beam dynamics in the storage ring
- Outstanding challenges for the future
- Summary and conclusions



Muon: (2nd generation lepton)

$\tau_{\mu} = 2.19703(4) \ \mu s$

$m_{\mu}c^2 = 105.658389(34) \text{ MeV}$

"Who ordered that?"



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Muon: (2nd generation lepton)

Source: weak decay $\pi \longrightarrow \mu$ $\overline{ u}_{\mu}$

The Pion Rest Frame





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Magnetic Moments: g-factors, etc

$$\vec{\mu}_s = g_s \left(\frac{e}{2m}\right) \vec{s}$$

 μ – magnetic moment g – g value s – spin

- Dirac Equation predicts $g \equiv 2$
- In nature, radiative corrections make g ≠ 2



An aside: The QED contribution to both electron and (g-2) is calculated through $(\alpha | \pi)^4$, with the $(\alpha | \pi)^5$ terr being calculated by Kinoshita et al.

 This level of precision is far beyond what its inver imagined.



"The main point was that all of us who put QED together, including especially Feynman, considered it a jerry-built and provisional structure which would either collapse or be replaced by something more permanent within a few years. So I find it amazing that it has lasted for fifty years and still agrees with experiments to twelve significant figures. It seems that Nature is telling us something. Perhaps she is telling us that she loves sloppiness." Freeman Dyson (private communication) – December 2006



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The SM value for lepton MDMs is one of the most precisely calculated numbers in physics

The Electron: to the level of the experimental error (4ppb), $a_e(\text{Standard Model}) = a_e(\text{QED with } \gamma, e)$

Contribution of μ , (or anything heavier than the electron) is ≤ 4 ppb.

For the muon, the relative contribution of heavier particles

 $\sim (rac{m_{\mu}}{m_{e}})^{2} \sim 40,000 \qquad \Rightarrow$

 $a_{\mu}(SM) = a_{\mu}(QED) + a_{\mu}(hadronic) + a_{\mu}(weak)$



Standard Model Value for (g-2)





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Standard Model Value for (g-2)





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The Dominant 2 π Channel



A. Hoecker Tau2010, Davier, Hoecker, Malaescu and Zhang, arXiv:1010.418v1 [hep-ph]



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$\Delta a_{\mu}(E821 - SM) = (296 \pm 81) \times 10^{-11}$

using e^+e^- data for the hadronic contribution.

An interesting, but not definitive difference with theory.



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Many BSM speculations: e.g SUSY (large tan β)

$$\frac{\tilde{\chi}}{\mu} \frac{\chi^{0}}{\chi^{-}} \frac{\chi^{0}}{\mu} + \frac{\chi^{0}}{\tilde{\mu}} \frac{\chi^{0}}{\chi^{-}} \frac{\mu}{\tilde{\mu}}$$

$$a_{\mu}(\text{SUSY}) \simeq \frac{\alpha(M_Z)}{8\pi \sin^2 \theta_W} \frac{m_{\mu}^2}{\tilde{m}^2} \tan \beta \left(1 - \frac{4\alpha}{\pi} \ln \frac{\tilde{m}}{m_{\mu}}\right)$$

 $\simeq (\text{sgn}\mu) \ 13 \times 10^{-10} \ \tan \beta \ \left(\frac{100 \text{ GeV}}{\tilde{m}}\right)^2$

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Other Models

- Technicolor
 - small Δa_{μ}
- Littlest Higgs with T-parity
 - small Δa_{μ}
- Universal Extra Dimensions
 - small Δa_{μ}
- Randall Sundrum
 - could accommodate large Δa_{μ}
- Two Higgs doublets, shadow Higgs
 - small Δa_{μ}
- Additional light bosons that can affect EM interactions (difficult to study at LHC)





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Traditionally,

• For many years, muon (g-2) has provided strong and serious constraints on models of physics beyond the standard model.





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Experimental Technique



Fermilab Beam Scheme



The 900-m long decay beam reduces the pion "flash" by x20 and leads to 6 – 12 times more stored muons per proton (compared to BNL)



Flash compared to BNL

parameter	FNAL/BNL
p / fill	0.25
π / p	0.4
π survive to ring	0.01
π at magic P	50
Net	0.05

Stored Muons / POT

parameter	BNL	FNAL	gain factor $\mathrm{FNAL}/\mathrm{BNL}$
\mathbf{Y}_{π} pion/p into channel acceptance	pprox 2.7E-5	$\approx 1.1\text{E-}5$	0.4
L decay channel length	88 m	$900 \mathrm{~m}$	2
decay angle in lab system	$3.8\pm0.5~\mathrm{mr}$	forward	3
$\delta p_{\pi}/p_{\pi}$ pion momentum band	$\pm 0.5\%$	$\pm 2\%$	1.33
FODO lattice spacing	$6.2 \mathrm{m}$	$3.25~\mathrm{m}$	1.8
inflector	closed end	open end	2
total			11.5

Plan View of the Injection Line





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The requirement of a uniform magnetic field means no gaps for injecting the beam. Everything has to fit between the pole pieces.



Inserting a Pole Piece



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Kapton insulation to prevent eddy currents from running around the ring, especially during an energy extraction or quench.

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muon (g-2) storage ring

H	Muon lifetime
-	(g-2) period
1	Cyclotron perio

of turns in ring

C

 $t_m = 64.4 ms$ $t_a = 4.37 ms$ $t_c = 149 ns$

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Mapping the Field



Free induction decay signals:





NMR B-field mapping trolley

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Mapping the Field



Fixed probes monitor dipole and quadrupole field components

Free induction decay signals:



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Azimuthal Variation



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Average Field Seen by the Muons $< B > = \int M(r, \theta) B(r, \theta) r dr d\theta$

In the average, B couples multipole by multipole with the moments of the muon distribution.

A circular aperture minimizes the effect of higher multipoles on $\langle B \rangle$

 $B(r,\theta) = \sum_{n=0} r^n \left[c_n \cos n\theta + s_n \sin n\theta \right]$

 ∞

 $M(r,\theta) = \sum [\gamma_m(r) \cos m\theta + \sigma_m(r) \sin m\theta]$



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$\langle B \rangle_{\phi}$ for 2001 averaged over azimuth 0.5 ppm contours









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The Kicker Modulator





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The Kicker Modulator





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Kicker Plate Geometry



electrodes



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The Kicker Current Pulse



B U eddy currents less than 0.1 ppm on B·dl after 20 μs



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The Electrostatic Quadrupoles: μ^+ polarity



 $\sim \pm 24$ kV at full power, 17 kV for beam scraping after injection



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The Ring Layout





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Scraping the Beam



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•
$$V_0 = \pm 24 \text{ kV}$$

- Beam is lifted and moved sideways
- Scraped on collimators to minimize losses

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Ring β -Function



 $x(s) = A\sqrt{\beta_x} \cos(\psi + \delta)$ $\psi(s) = \sqrt{Ks}$

 $\frac{\beta_{max}}{2} = 1.03$ min

for 4-fold symmetry



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Weak Focusing Betatron

Field index: $n = \frac{\kappa R_0}{\beta B_0} \simeq 0.135$ $f_y = f_C \sqrt{n} \simeq 0.37 f_C;$

$$f_x = f_C \sqrt{1 - n} \simeq 0.929 f_C$$

 Detector acceptance depends on the radial coordinate x. The beam moves coherently radially relative to a detector with the "Coherent Betatron Frequency" (CBO)

$$f_{CBO} = f_C - f_x = (1 - \sqrt{1 - n})f_C$$





CBO amplitude modulates the signal in the detectors.



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Tune Plane VV $2v_y = 1$ 0.50 $\sqrt{1}$ – ν_x = \boldsymbol{n} 0.45 21 × $\nu_y =$ 0.40 n n = 0.148 0.35 n = 0.1000.30 50 0.25 $V_{\mathbf{X}}$ 1.00 0.80 0.85 0.90 0.95



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Muon Decay

 μ -decay: parity violating $\mu^- \rightarrow e^- \overline{\nu}_e \nu_\mu$

The Muon Rest Frame





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e^{\pm} from $\mu^{\pm} \rightarrow e \pm \nu \nu$ are detected





400 MHz digitizer gives t, E

Count number of e⁻ with *E_e* ≥ 1.8 GeV



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4×10^9 e^- , $E_{e^-} \ge$ 1.8 GeV $= N_0 e^{-\lambda t} [1 + A \cos \omega_a t + \phi)]$

electron time spectrum (2001)



In the 1999 Data Set: A Surprise Nature gives us 5 parameters: $f(t) = N_0 e^{-\lambda t} [1 + A \cos(\omega_a t + \phi)]$

Storage ring plus bunched beam gives us more:

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Frequencies in the (g-2) Ring

Quantity	Expression	Frequency	Period	
f_a	$rac{e}{2\pi mc}a_{\mu}B$	0.23 MHz	4.37 μs	
f_c	$rac{v}{2\pi R_0}$	6.7 MHz	149 ns	
f_x	$\sqrt{1-n}f_c$	6.23 MHz	160 ns	
f_y	$\sqrt{n}f_c$	2.48 MHz	402 ns	
$f_{\rm CBO}$	$f_c - f_x$	0.477 MHz	2.10 μ s	
$f_{ m VW}$	$f_c - 2f_y$	1.74 MHz	$0.574 \ \mu s$	



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Fiber Beam Monitors





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Measuring the Tune from a single fiber



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The Tune During Scraping



• The tune change with scraping is clearly visible from the fiber harps



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CBO in the 2001 Data Set $f(t) = N_0 e^{-\lambda t} [1 + A \cos(\omega_a t + \phi)]$

Residuals from fitting the 5-parameter function



Beam Debunching after Injection



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What do we expect for distribution of equilibrium radii?



Equilibrium orbit distribution

Distribution of Equilibrium Radii



E989: Systematic Error Goal

Systematic uncertainty (ppm)	1998	1999	2000	2001	E989 Goal
Magnetic field – ω_p	0.5	0.4	0.24	0.17	0.07
Anomalous precession – ω_a	0.8	0.3	0.3	0.21	0.07



Systematic errors on ω_a (ppm)



 $\Sigma^{*} = 0.11$



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Outstanding Beam Issues

- The Kicker: can it be improved?
 - What's the real injection efficiency in E821 and can it be simulated?
- Elimination of the CBO



We have developed a new simulation tool to guide our improvements. It includes all major subsystems



- Vacuum vessel
- Inflector
- Kickers
- Quadrupoles
- Collimators



Example: Collective beam motion for stored muons is reproduced



The predicted horizontal and vertical betatron oscillation frequencies match measured values from E821 at the percent level.



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We can simulate modified kicker pulse shapes to predict storage improvements



in the NIM paper we said:

but ...



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Relative Amplitude of the CBO effect if not accounted for



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Lost Muons and CBO are Major Issues

- Two schemes proposed to eliminate CBO and losses
 - Drive CBO with an oscillating dipole to scrape, then slip the phase by π and damp it
 - Suggested by Yuri Orlov
 - Pulsed Octupole for 30 turns
 - Suggested by Yuri Shatunov



Oscillating Dipole Solution



 Use Fiber Harps to measure phase of CBO

Sample Parameters

L = 0.5 m N = 20 turns E_{x0} = 7.4 kV/cm f = 470 kHz



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CBO cure: betatron phases mixing by nonlinear fields

2-meter long coil!

$$\Delta B_{x} + i\Delta B_{y} = B_{0} \sum_{n} (b_{n} + ia_{n})(x + iy)^{n}$$

$$octupole: (n = 3)$$

$$\beta \left(\frac{\partial v}{\partial a^{2}}\right) \approx \frac{3}{8} \langle \beta^{2} b_{3} \rangle_{s}$$

$$\int \frac{1}{\sqrt{2}} \int \frac{1}{\sqrt{2}} \frac{3}{8} \langle \beta^{2} b_{3} \rangle_{s}$$

$$\int \frac{1}{\sqrt{2}} \int \frac{1}{\sqrt{2}} \frac{3}{8} \langle \beta^{2} b_{3} \rangle_{s}$$

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$$\int \frac{1}{\sqrt{2}} \int \frac{1}{\sqrt{2}} \frac{1}{\sqrt{2}} \frac{3}{8} \langle \beta^{2} b_{3} \rangle_{s}$$

$$\int \frac{1}{\sqrt{2}} \int \frac{1}{\sqrt{2}} \frac{1}{\sqrt$$

Y. Shatunov, SPIN04



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Challenges with Octupole

- Eddy currents affecting B₀?
 - We can only tolerate effects on $B \cdot dl$ at the < 0.03 ppm level
- Too many muons lost?



Summary

- E821 at BNL achieved 0.54 ppm relative accuracy on a_{μ}
 - -0.46 ppm statistical
 - 0.28 ppm systematic
- This represents a factor of 14 over the CERN experiment



Summary

- E821 at BNL achieved 0.54 ppm relative accuracy
 - 0.46 ppm statistical
 - 0.28 ppm systematic
- This represents a factor of 14 over the CERN experiment
- P/E989 Aims to achieve an additional factor of 4
 from 0.54 ppm → 0.14 ppm
- Will more than double the physics reach when confronting theory
- Many kind words, no \$ yet.

-Please come join us!





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Traceback: π Injection vs. μ Injection







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a(had) from hadronic τ decay? e^{+} γ h τ^{-} v_{τ} e^{+} γ π^{-} v_{τ} W^{-} π^{-} e^{-} γ h e^{-} γ h τ^{-} v_{τ} ψ^{-} ψ^{-

- Assume: CVC, no 2nd-class currents, isospin breaking corrections.
- n.b. τ decay has no isoscalar piece, while e⁺e⁻ does
- Many inconsistencies in comparison of e⁺e⁻ and
 - τ decay:

- Using e^+e^- data and CVC to predict τ branching ratio gives

2.1 to 3.6 σ discrepancies with reality.

- F_{π} from τ decay has different shape from e⁺e⁻.

In summary, there are unanswered questions from the τ data.



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Where we came from:



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Today with e^+e^- based theory:



 $a_{\mu} = 116592089(63) \times 10^{-11} (0.54 \text{ ppm})$



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Magnetic Circuits



 $\Phi \oint \frac{d\ell}{\mu A} = NI \quad \Phi \Re = MMF$ Ohm's law



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Schematic of the Magnet





Winding the Coils





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The Finished Coils





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Coil Interconnect





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