

F → SP(2N) [for notes]



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Generating Functions produce symplectic transport maps

$$F_1(\vec{q}, \vec{q}_0, s) \quad \text{with} \quad \vec{p} = -\vec{\partial}_q F_1(\vec{q}, \vec{q}_0, s) \quad , \quad \vec{p}_0 = \vec{\partial}_{q_0} F_1(\vec{q}, \vec{q}_0, s)$$

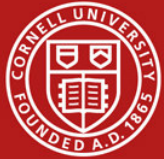
$$\left. \begin{aligned} \vec{z} &= \begin{pmatrix} \vec{q} \\ \vec{p} \end{pmatrix} = \begin{pmatrix} \vec{q} \\ -\vec{\partial}_q F_1(\vec{q}, \vec{q}_0, s) \end{pmatrix} = \vec{f}(\vec{Q}, s) \\ \vec{z}_0 &= \begin{pmatrix} \vec{q}_0 \\ \vec{p}_0 \end{pmatrix} = \begin{pmatrix} \vec{q}_0 \\ \vec{\partial}_{q_0} F_1(\vec{q}, \vec{q}_0, s) \end{pmatrix} = \vec{g}(\vec{Q}, s) \end{aligned} \right\} \begin{aligned} \vec{z} &= \vec{f}(\vec{g}^{-1}(\vec{z}_0, s), s) \\ \vec{M} &= \vec{f} \circ \vec{g}^{-1} \\ &\text{(function concatenation)} \end{aligned}$$

Jacobi matrix of concatenated functions:

$$\vec{C}(\vec{z}_0) = \vec{A} \circ \vec{B}(\vec{z}_0)$$

$$C_{ij} = \partial_j C_i = \sum_k \partial_{z_{0j}} B_k(\vec{z}_0) \left[\partial_{z_k} A_i(\vec{z}) \right]_{\vec{z}=\vec{B}(\vec{z}_0)} \Rightarrow \underline{C} = \underline{A}(\underline{B})\underline{B}$$

$$\vec{M} \circ \vec{g} = \vec{f} \Rightarrow \underline{M}(\underline{g}) = \underline{F}\underline{G}^{-1}$$



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$$\vec{f}(\vec{Q}, s) = \begin{pmatrix} \vec{q} \\ -\vec{\partial}_q F_1(\vec{q}, \vec{q}_0, s) \end{pmatrix} \Rightarrow F = \begin{pmatrix} 1 & 0 \\ -\vec{\partial}_q \vec{\partial}_q^T F_1 & -\vec{\partial}_q \vec{\partial}_{q_0}^T F_1 \end{pmatrix}$$

$$\vec{g}(\vec{Q}, s) = \begin{pmatrix} \vec{q}_0 \\ \vec{\partial}_{q_0} F_1(\vec{q}, \vec{q}_0, s) \end{pmatrix} \Rightarrow G = \begin{pmatrix} 0 & 1 \\ \vec{\partial}_{q_0} \vec{\partial}_q^T F_1 & \vec{\partial}_{q_0} \vec{\partial}_{q_0}^T F_1 \end{pmatrix}$$

$$F = \begin{pmatrix} 1 & 0 \\ -F_{11} & -F_{12} \end{pmatrix}, \quad G = \begin{pmatrix} 0 & 1 \\ F_{21} & F_{22} \end{pmatrix} \Rightarrow G^{-1} = \begin{pmatrix} -F_{21}^{-1} F_{22} & F_{21}^{-1} \\ 1 & 0 \end{pmatrix}$$

$$\underline{M}(\vec{g}) = FG^{-1} = \begin{pmatrix} -F_{21}^{-1} F_{22} & F_{21}^{-1} \\ F_{11} F_{21}^{-1} F_{22} - F_{12} & -F_{11} F_{21}^{-1} \end{pmatrix}$$

$$\underline{M} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \underline{M}^T \longrightarrow$$

The map from a generating function is symplectic.

$$= \begin{pmatrix} -F_{21}^{-1} & -F_{21}^{-1} F_{22} \\ F_{11} F_{21}^{-1} & F_{11} F_{21}^{-1} F_{22} - F_{12} \end{pmatrix} \begin{pmatrix} -F_{22} F_{12}^{-1} & F_{22} F_{12}^{-1} F_{11} - F_{21} \\ F_{12}^{-1} & -F_{12}^{-1} F_{11} \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$



Symplectic transport maps have a Generating Functions

$$\vec{z} = \vec{M}(\vec{z}_0)$$

$$\begin{pmatrix} \vec{q} \\ \vec{q}_0 \end{pmatrix} = \begin{pmatrix} \vec{M}_1(\vec{z}_0) \\ \vec{q}_0 \end{pmatrix} = \vec{l}(\vec{z}_0), \quad \begin{pmatrix} \vec{p}_0 \\ \vec{p} \end{pmatrix} = \begin{pmatrix} \vec{p}_0 \\ \vec{M}_2(\vec{z}_0) \end{pmatrix} = \vec{h}(\vec{z}_0) = \underline{J} \left[\vec{\partial} F_1(\vec{q}, \vec{q}_0) \right]_{\vec{l}(\vec{z}_0)}$$

$$\vec{\partial} F_1 = -\underline{J} \vec{h} \circ \vec{l}^{-1} = \vec{F}$$

For F_1 to exist it is necessary and sufficient that $\partial_i F_j = \partial_j F_i \Rightarrow \underline{F} = \underline{F}^T$

$$-\underline{J} \vec{h} = \vec{F} \circ \vec{l} \Rightarrow -\underline{J} \underline{h} = \underline{F}(\vec{l}) \underline{l}$$

Is $\underline{J} \underline{h} \underline{l}^{-1}$ symmetric? Yes since:

$$\begin{aligned} \underline{J} \underline{h} \underline{l}^{-1} &= \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ \vec{\partial}_{q_0}^T \vec{M}_2 & \vec{\partial}_{p_0}^T \vec{M}_2 \end{pmatrix} \begin{pmatrix} \vec{\partial}_{q_0}^T \vec{M}_1 & \vec{\partial}_{p_0}^T \vec{M}_1 \\ 1 & 0 \end{pmatrix}^{-1} \\ &= \begin{pmatrix} M_{21} & M_{22} \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ M_{12}^{-1} & -M_{12}^{-1} M_{11} \end{pmatrix} = \begin{pmatrix} M_{22} M_{12}^{-1} & M_{21} - M_{22} M_{12}^{-1} M_{11} \\ M_{12}^{-1} & M_{12}^{-1} M_{11} \end{pmatrix} \end{aligned}$$



SP(2N) → F [for notes]



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$$\underline{Jh} \underline{l}^{-1} = \begin{pmatrix} M_{22} M_{12}^{-1} & M_{21} - M_{22} M_{12}^{-1} M_{11} \\ M_{12}^{-1} & M_{12}^{-1} M_{11} \end{pmatrix} = \begin{pmatrix} A & B \\ C & D \end{pmatrix}$$

$$\vec{M}(\vec{z}_0) = \begin{pmatrix} \vec{M}_1(\vec{q}_0, \vec{p}_0) \\ \vec{M}_2(\vec{q}_0, \vec{p}_0) \end{pmatrix}, \quad \underline{M} = \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix}$$

$$\underline{M} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \underline{M}^T = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

$$\begin{pmatrix} -M_{12} & M_{11} \\ -M_{22} & M_{21} \end{pmatrix} \begin{pmatrix} M_{11}^T & M_{21}^T \\ M_{12}^T & M_{22}^T \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

$$M_{12} M_{11}^T = M_{11} M_{12}^T \quad \Rightarrow \quad (M_{12}^{-1} M_{11})^T = [M_{12}^{-1} M_{11} M_{12}^T] M_{12}^{-T} = M_{12}^{-1} M_{11}$$

$$M_{21} M_{22}^T = M_{22} M_{21}^T$$

$$M_{11} M_{22}^T - M_{12} M_{21}^T = 1$$

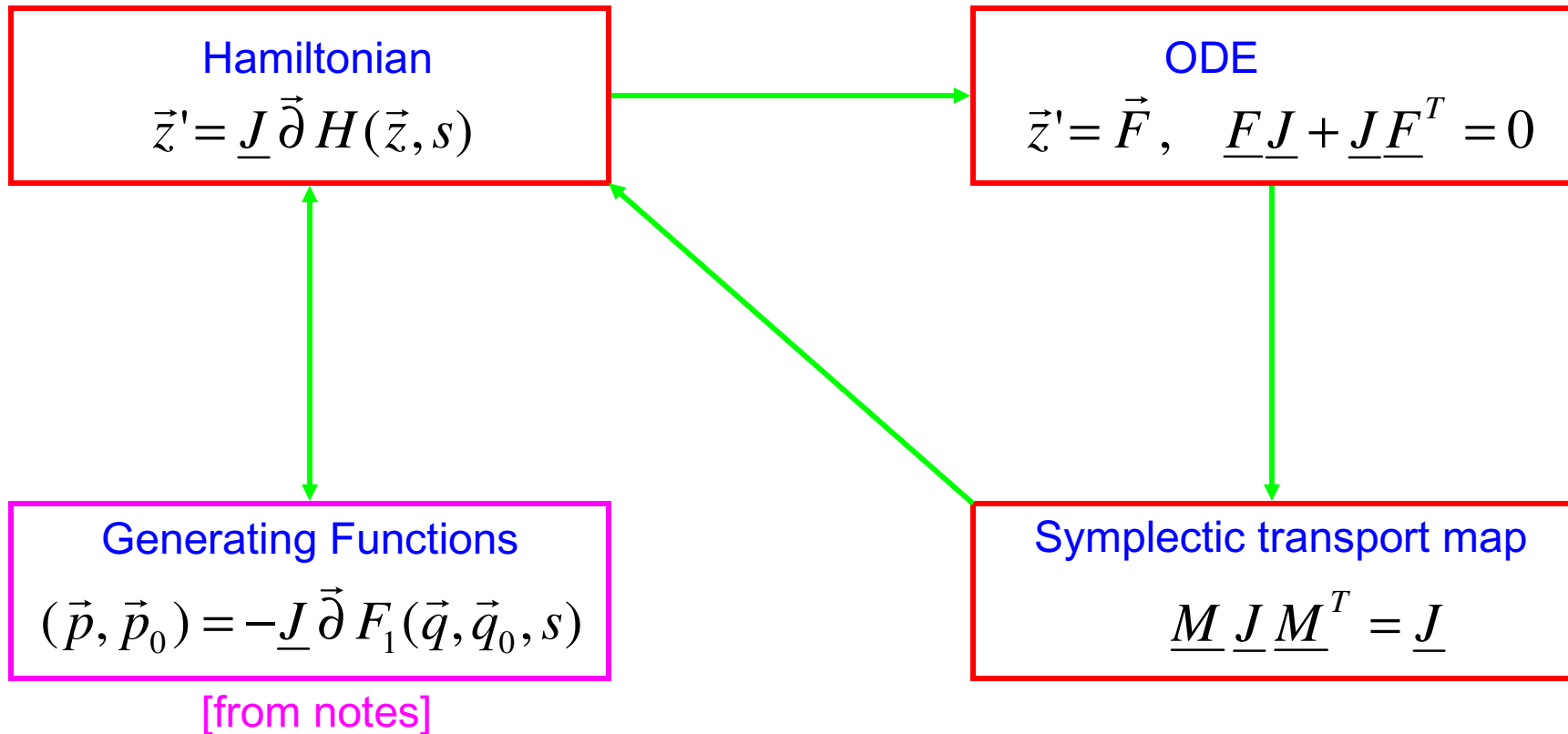
$$M_{22} M_{11}^T - M_{21} M_{12}^T = 1$$

$$D = D^T$$

$$A = A^T$$

$$(M_{22} M_{12}^{-1})^T = [M_{22} M_{11}^T M_{12}^{-T} - M_{21}] M_{22}^T = M_{22} [M_{12}^{-1} M_{11} M_{22}^T - M_{21}^T] = M_{22} M_{12}^{-1}$$

$$M_{21} - M_{22} M_{12}^{-1} M_{11} = M_{21} - M_{22} M_{11} M_{12}^{-T} = M_{12}^{-T} \quad \longrightarrow \quad B = C^T$$





Advantages of Symplecticity



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- Determinant of the transfer matrix of linear motion is 1:

$$\vec{z}(s) = \underline{M}(s) \cdot \vec{z}_0 \quad \text{with} \quad \det(\underline{M}(s)) = +1$$

- One function suffices to compute the total nonlinear transfer map:

$$F_1(\vec{q}, \vec{q}_0, s) \quad \text{with} \quad \vec{p} = -\vec{\partial}_q F_1(\vec{q}, \vec{q}_0, s) \quad , \quad \vec{p}_0 = \vec{\partial}_{q_0} F_1(\vec{q}, \vec{q}_0, s)$$

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- Therefore Taylor Expansion coefficients of the transport map are related.
- Computer codes can numerically approximate $\vec{M}(s, \vec{z}_0)$ with exact symplectic symmetry.
- Liouville's Theorem for phase space densities holds.



Eigenvalues of a Symplectic Matrix



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For matrices with real coefficients:

If there is an eigenvector and eigenvalue: $\underline{M}\vec{v}_i = \lambda_i \vec{v}_i$

then the complex conjugates are also eigenvector and eigenvalue: $\underline{M}\vec{v}_i^* = \lambda_i^* \vec{v}_i^*$

For symplectic matrices:

If there are eigenvectors and eigenvalues: $\underline{M}\vec{v}_i = \lambda_i \vec{v}_i$ with $\underline{J} = \underline{M}^T \underline{J} \underline{M}$

then $\vec{v}_i^T \underline{J} \vec{v}_j = \vec{v}_i^T \underline{M}^T \underline{J} \underline{M} \vec{v}_j = \lambda_i \lambda_j \vec{v}_i^T \underline{J} \vec{v}_j \Rightarrow \vec{v}_i^T \underline{J} \vec{v}_j (\lambda_i \lambda_j - 1) = 0$

Therefore $\underline{J} \vec{v}_j$ is orthogonal to all eigenvectors with eigenvalues that are not $1/\lambda_j$. Since it cannot be orthogonal to all eigenvectors, there is at least one eigenvector with eigenvalue $1/\lambda_j$

Two dimensions: λ_j is eigenvalue

Then $1/\lambda_j$ and λ_j^* are eigenvalues

$$\lambda_2 = 1/\lambda_1 = \lambda_1^* \Rightarrow |\lambda_j| = 1$$

$$\lambda_2 = 1/\lambda_1 = \lambda_1^*$$

