



$$\frac{d}{d\vartheta} J_x = \cos(\tilde{\psi}_x + \varphi_x) \sqrt{2J_x \beta_x} \Delta f_x \frac{L}{2\pi} \quad , \quad \frac{d}{d\vartheta} \varphi_x = \nu_x - \sin(\tilde{\psi}_x + \varphi_x) \sqrt{\frac{\beta_x}{2J_x}} \Delta f_x \frac{L}{2\pi}$$

$$\frac{d}{d\vartheta} J_y = \cos(\tilde{\psi}_y + \varphi_y) \sqrt{2J_y \beta_y} \Delta f_y \frac{L}{2\pi} \quad , \quad \frac{d}{d\vartheta} \varphi_y = \nu_y - \sin(\tilde{\psi}_y + \varphi_y) \sqrt{\frac{\beta_y}{2J_y}} \Delta f_y \frac{L}{2\pi}$$

$$\frac{d}{d\vartheta} \vec{\varphi} = \vec{\partial}_J H \quad , \quad \frac{d}{d\vartheta} \vec{J} = -\vec{\partial}_\varphi H \quad , \quad H(\vec{\varphi}, \vec{J}, \vartheta) = \vec{\nu} \cdot \vec{J} - \frac{L}{2\pi} \int_0^{\vec{x}} \Delta \vec{f}(\hat{x}, s) d\hat{x}$$

The integral form can be chosen since it is path independent. This is due to the Hamiltonian nature of the force:

$$\Delta f_{x,y}(x, y, s) = -\partial_{x,y} \Delta H(x, y, s)$$

Single Resonance model for two dimensions means retaining only the amplitude dependent tune shift and one term in the two dimensional Fourier expansion:

$$H(\vec{\varphi}, \vec{J}, \vartheta) = \vec{\nu} \cdot \vec{J} + H_{00}(\vec{J}) + H_{n\vec{m}}(\vec{J}) \cos(n\vartheta + m_x \varphi_x + m_y \varphi_y + \Psi_{n\vec{m}}(\vec{J}))$$

For $n + m_x \nu_x + m_y \nu_y \approx 0$

$$m_x \varphi_x + m_y \varphi_y = \vec{m} \cdot \vec{\varphi}$$



$n + m_x \nu_x + m_y \nu_y \approx 0$ means that oscillations in y can drive oscillations in x in

$$x'' = -K x + \Delta f_x(x, y, s)$$

The resonance term in the Hamiltonian then changes only slowly:

$$H(\vec{\varphi}, \vec{J}, \vartheta) = \vec{v} \cdot \vec{J} + H_{00}(\vec{J}) + H_{n\vec{m}}(\vec{J}) \cos(n\vartheta + \vec{m} \cdot \vec{\varphi} + \Psi_{n\vec{m}}(\vec{J}))$$

$$\frac{d}{d\vartheta} \vec{\varphi} = \vec{\partial}_{\vec{J}} H \quad , \quad \frac{d}{d\vartheta} \vec{J} = -\vec{\partial}_{\varphi} H$$

$$J = \vec{m} \cdot \vec{J}$$

$$J_{\perp} = m_x J_x - m_y J_y = \vec{m} \times \vec{J} \quad \Rightarrow \quad \frac{d}{d\vartheta} J_{\perp} = 0$$

Difference resonances lead to stable motion since:

$$n + |m_x| \nu_x - |m_y| \nu_y \approx 0 \Rightarrow |m_x| J_x + |m_y| J_y = \text{const.}$$

Sum resonances lead to unstable motion since:

$$n + |m_x| \nu_x + |m_y| \nu_y \approx 0 \Rightarrow |m_x| J_x - |m_y| J_y = \text{const.}$$



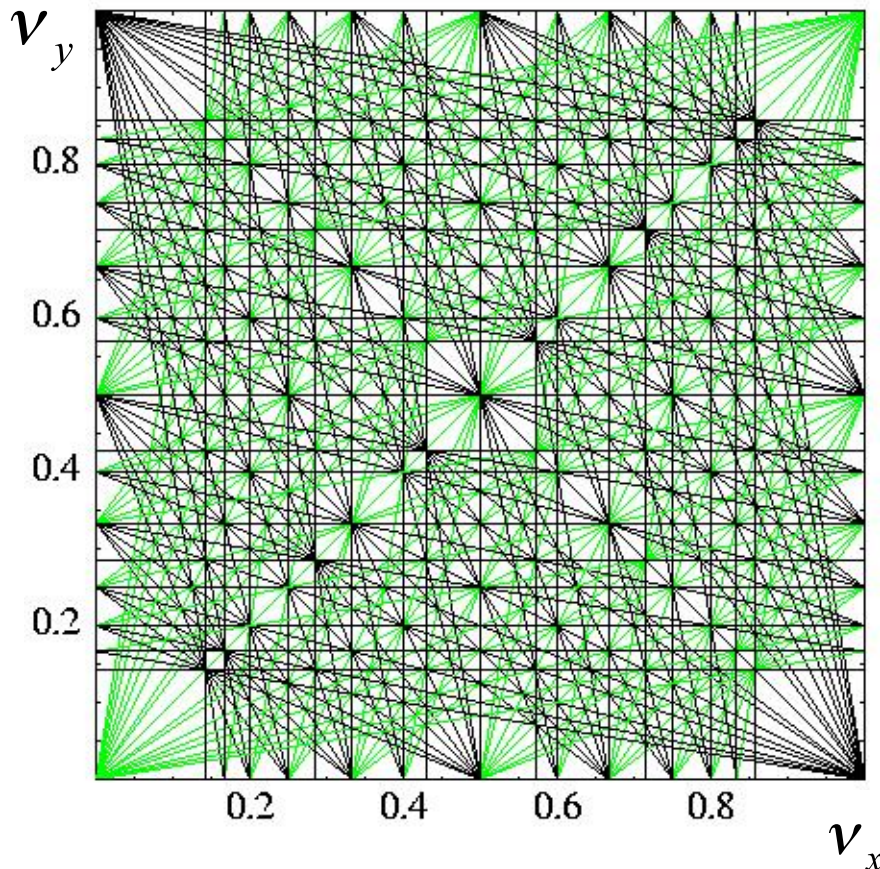
Resonances Diagram



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$n + m_x \nu_x + m_y \nu_y \approx 0$ means that oscillations in y can drive oscillations in x in

$$x'' = -K x + \Delta f_x(x, y, s)$$



All these resonances have to be avoided by their respective resonance width.

The position of an accelerator in the tune plane is called its Working Point.