



Diffraction (contd.)

Ivan Bazarov

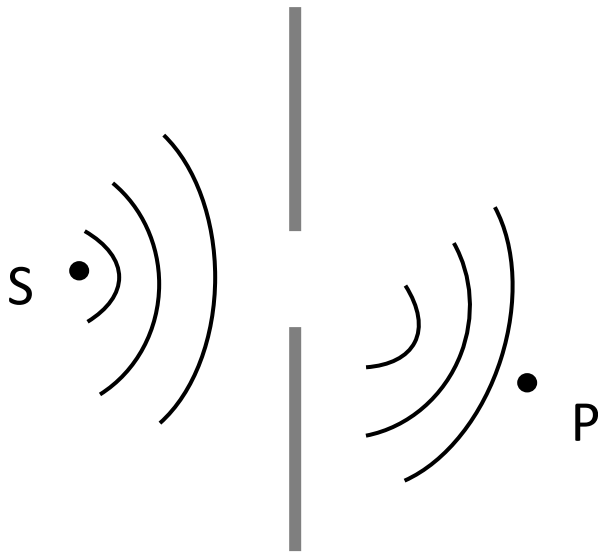
Cornell Physics Department / CLASSE

Outline

- **Fresnel vs. Fraunhofer diffraction**
- **Talbot effect**
- **Fresnel zones & plates**
- **Fraunhofer diffraction**
- **2D Fourier Transforms**

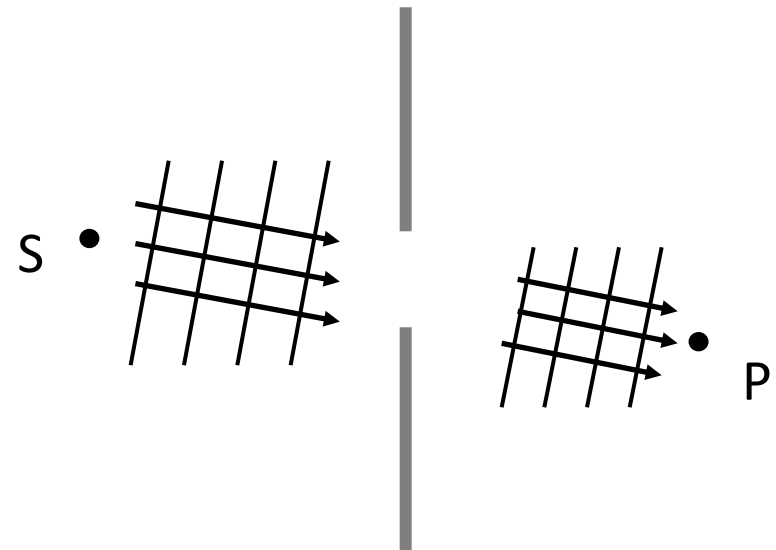


Fresnel vs. Fraunhofer diffraction



Fresnel:

occurs when either S or P are close enough to the aperture that wavefront curvature is not negligible



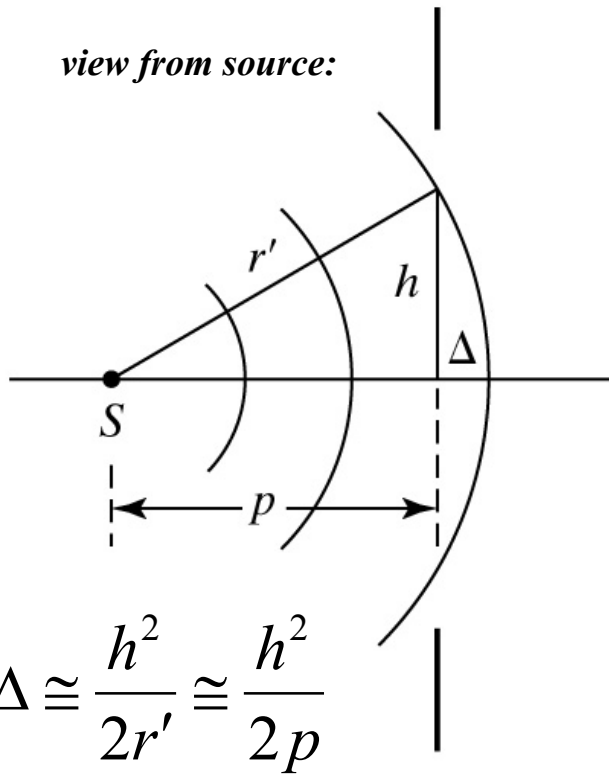
Fraunhofer:

both incident and diffracted waves may be considered to be planar (i.e. both S and P are far from the aperture)



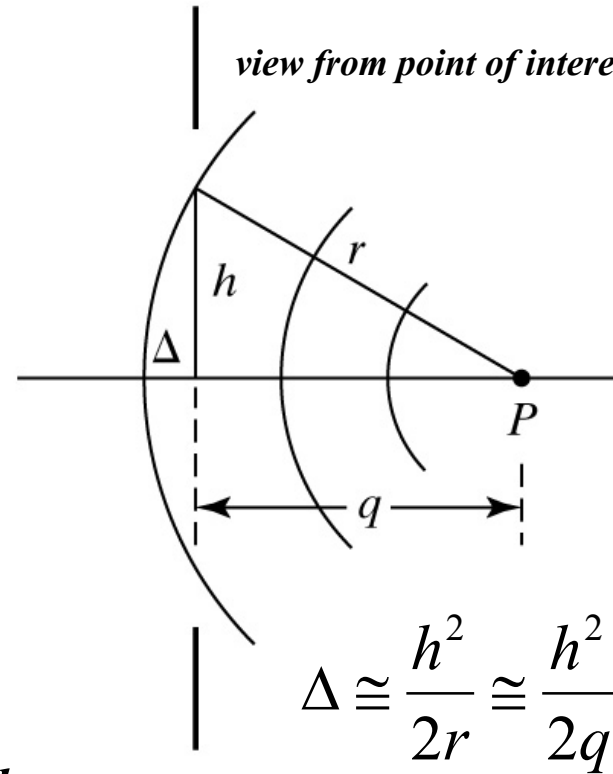
Fresnel vs. Fraunhofer criterion

view from source:



$$\Delta \cong \frac{h^2}{2r'} \cong \frac{h^2}{2p}$$

view from point of interest:



$$\Delta \cong \frac{h^2}{2r} \cong \frac{h^2}{2q}$$

near field \equiv

$$\Delta > \lambda$$

$$\frac{1}{2} \left(\frac{1}{p} + \frac{1}{q} \right) h^2 > \lambda$$

$$d < \frac{A}{\lambda}$$

where d represents p or q (=distance from source or point to aperture)

A is aperture area



Fresnel number

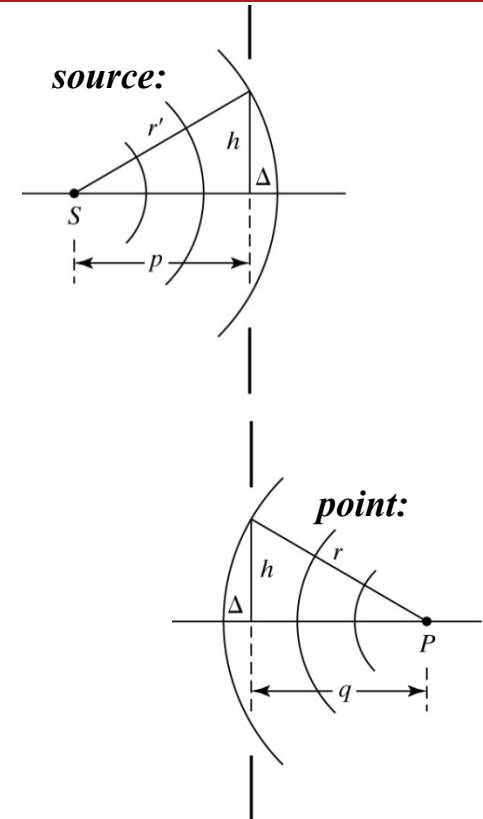
Fraunhofer diffraction occurs when:

$$F = \frac{h^2}{d\lambda} \ll 1$$

Fresnel diffraction occurs when:

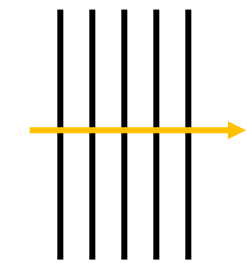
$$F = \frac{h^2}{d\lambda} \geq 1$$

where h = aperture or slit size
 λ = wavelength
 d = distance from the aperture (p or q)

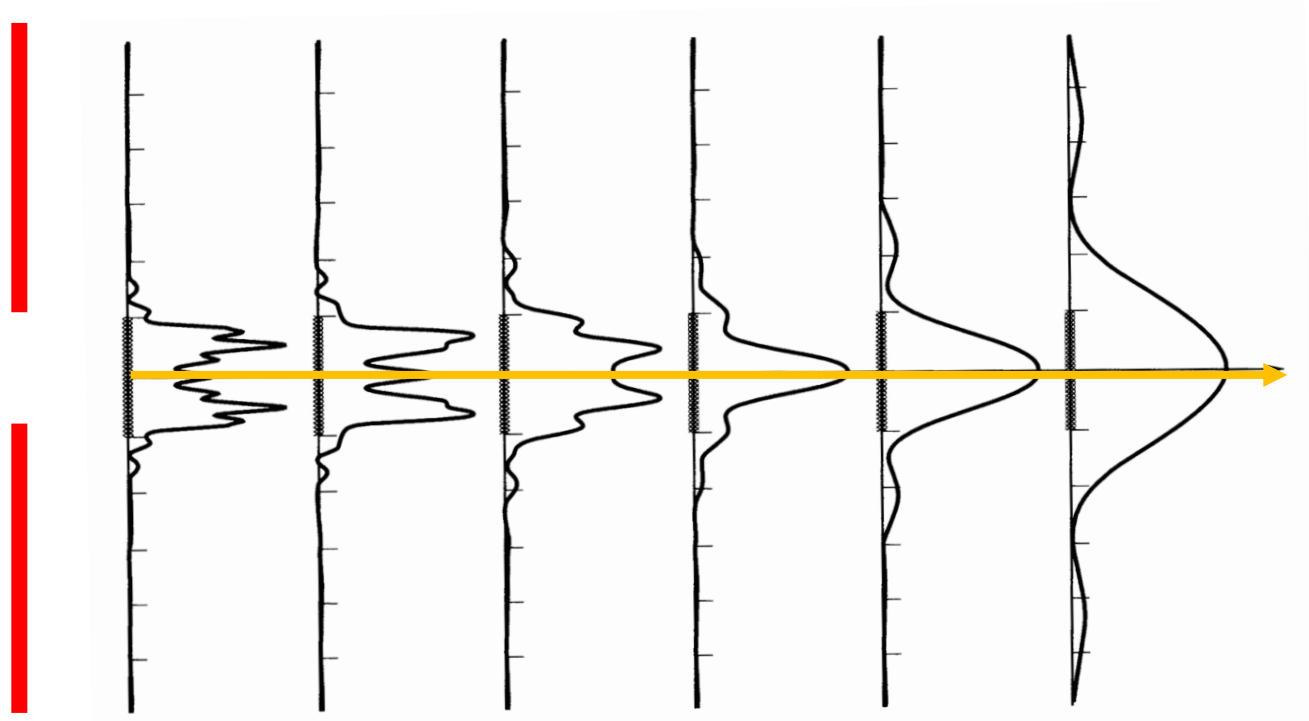




From Fresnel to Fraunhofer diffraction



Incident
plane
wave

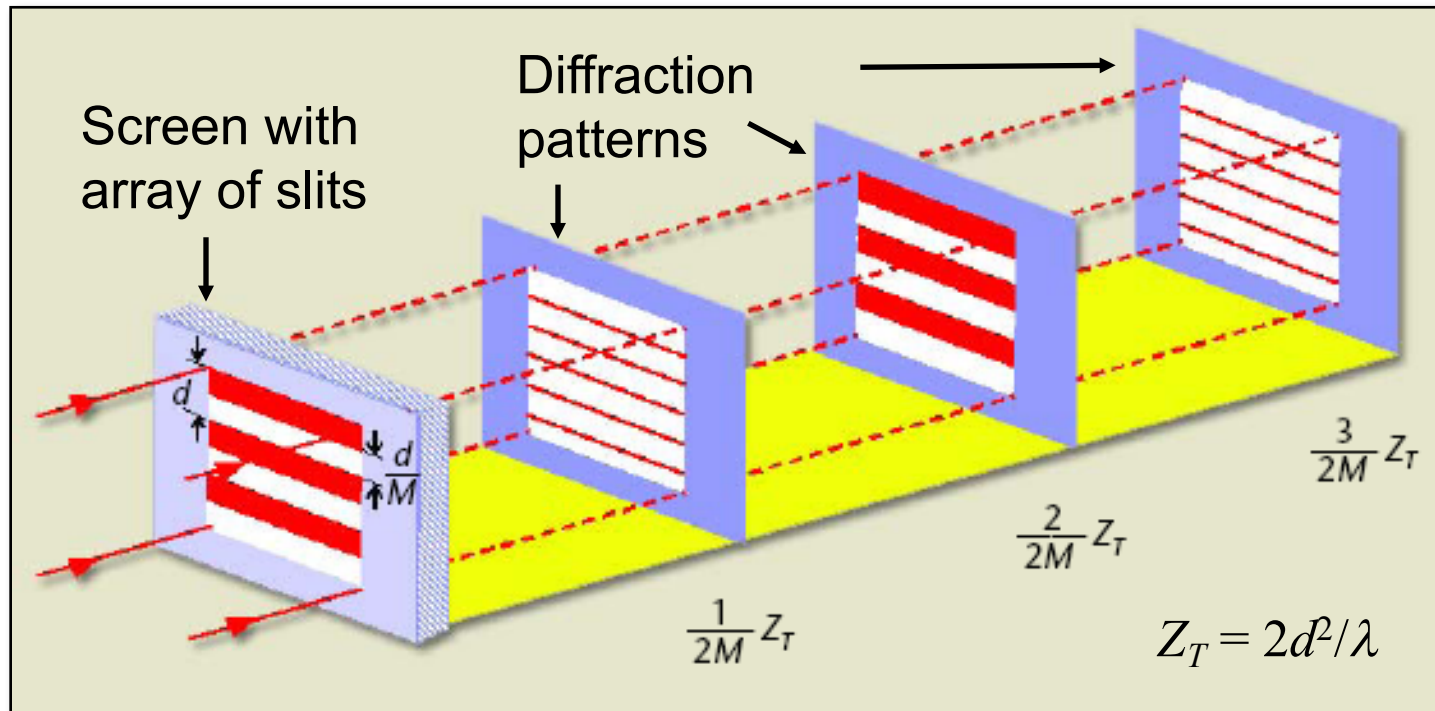


$F \gg 1$

$F \ll 1$



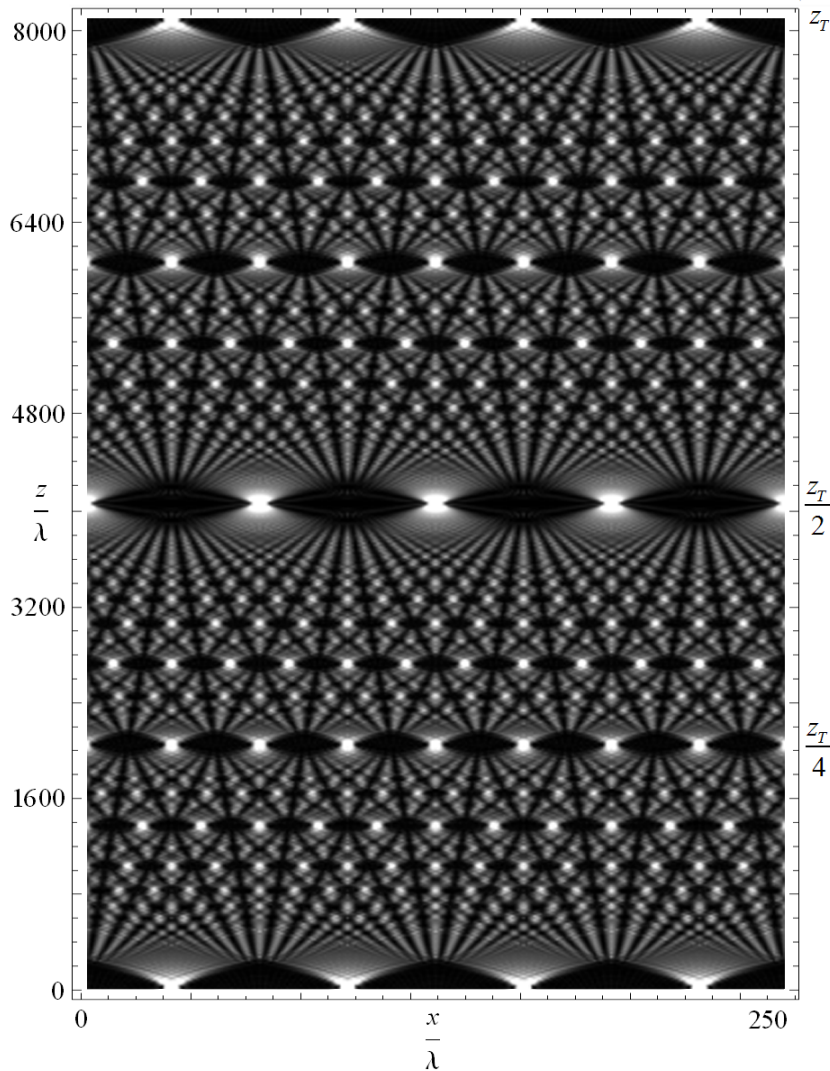
Fresnel diffraction from infinite array of slits: Talbot effect



- one of the few Fresnel diffraction problems that can be solved analytically
- beam pattern alternates between two different fringe patterns



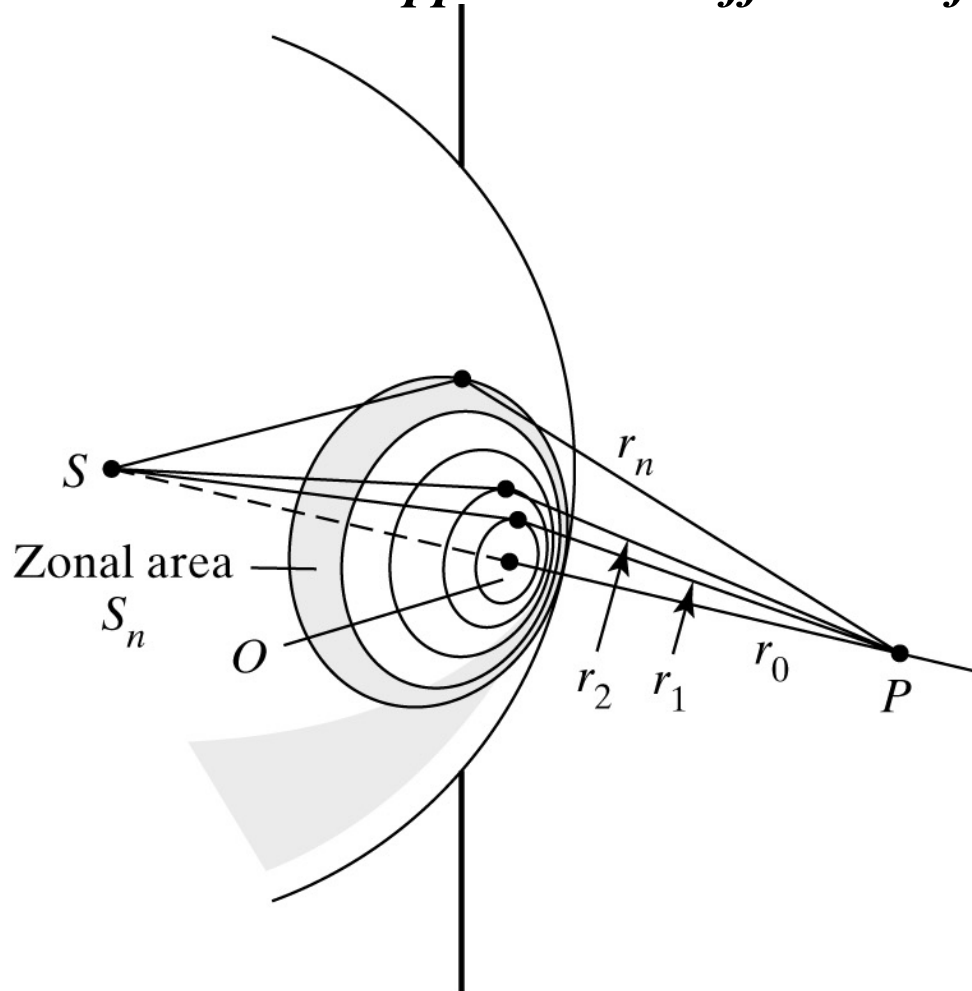
Talbot "carpet"





Fresnel zones (180° phase difference)

Fresnel's approach to diffraction from circular apertures



zone spacing = $\lambda/2$:

$$r_1 = r_0 + \lambda/2$$

$$r_2 = r_0 + \lambda$$

$$r_3 = r_0 + 3\lambda/2$$

\vdots

$$r_n = r_0 + n\lambda/2$$

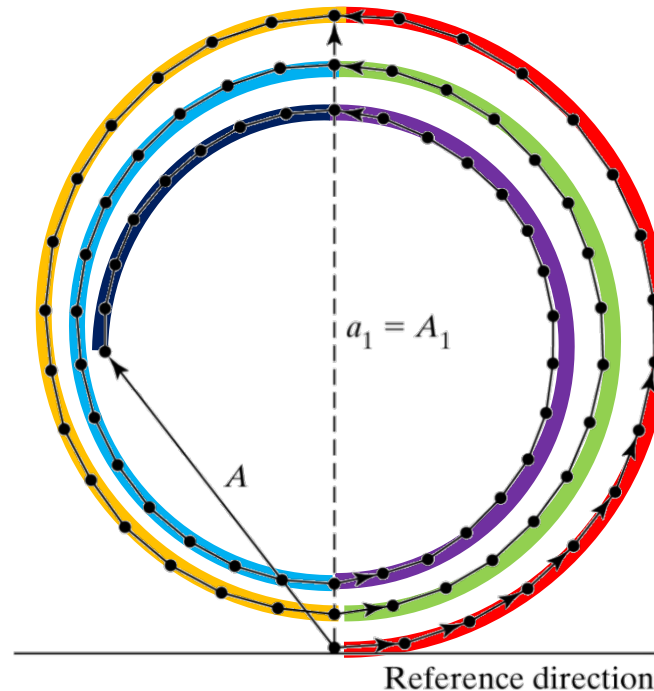
**these are called the
*Fresnel zones***

**(note: all zones
have equal areas)**



Adding up light from the zones

as we draw a phasor diagram where *each zone is subdivided into 15 subzones*



5¹/₂ half-period zones

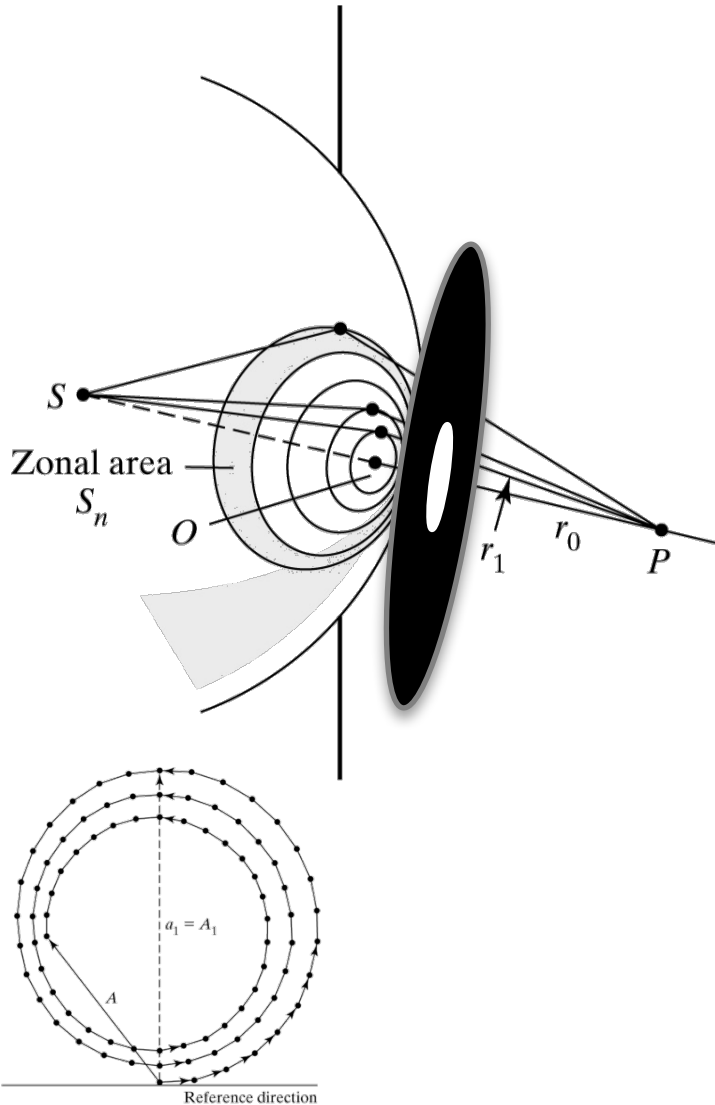
- obliquity factor shortens successive phasors
- circles do not close, but spiral inwards
- amplitude $a_1 = A_1$: resultant of subzones in 1st half-period zone
- composite amplitude at P from n half-period zones:

$$A_n = a_1 + a_2 e^{i\pi} + a_3 e^{i2\pi} + a_4 e^{i3\pi} + \dots a_n e^{i(n-1)\pi}$$

$$A_n = a_1 - a_2 + a_3 - a_4 + \dots a_n$$



Some interesting implications of Fresnel zones



A circular aperture is matched in size with the first Fresnel zone:

What is amplitude of the wavefront at P?

$$A_P = a_1$$

Now open the aperture wider to also admit zone 2:

$$A_P \sim 0 !$$

Now remove aperture, allowing all zones to contribute:

$$A_P = \frac{1}{2} a_1 !!!$$

(To find intensity – square the amplitudes, i.e. it's only $\frac{1}{4}$ of the 1st zone!)



Some interesting implications of Fresnel zones

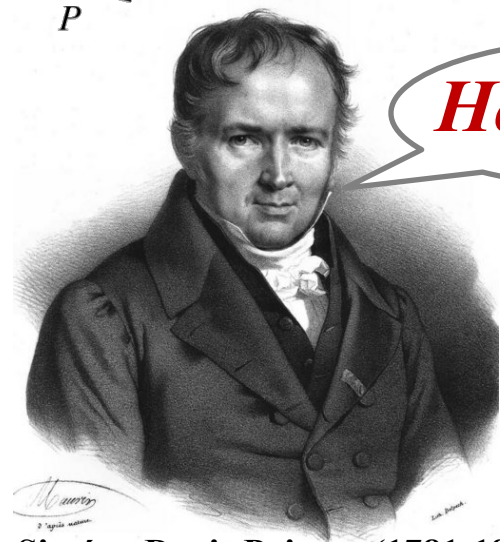
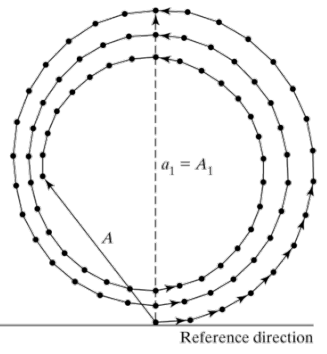
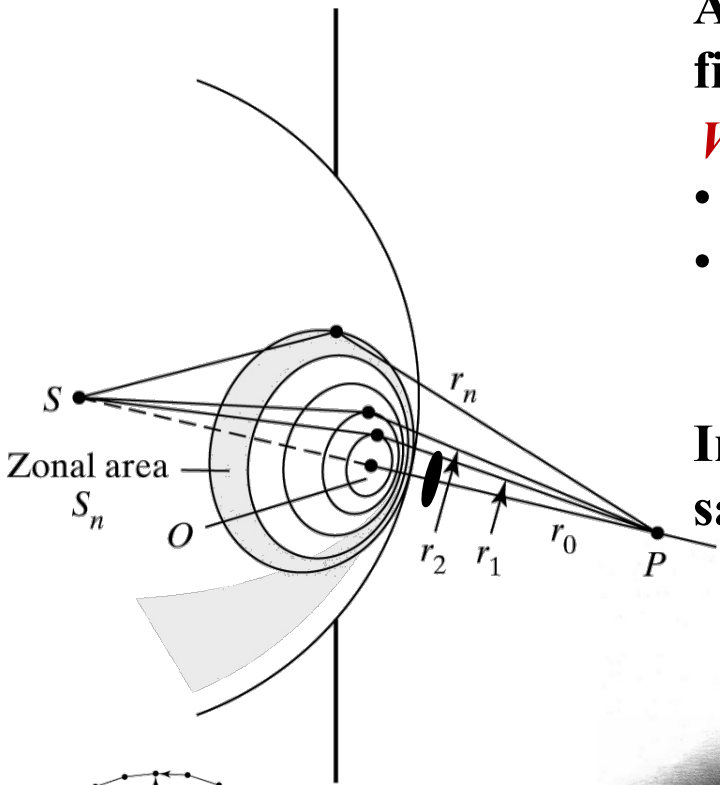
A circular disk is matched in size with the first Fresnel zone:

What is amplitude of the wavefront at P?

- all zones *except* the first contribute
- first contributing zone is the second

$$A_P = \frac{1}{2} a_2$$

Irradiance at center of shadow nearly the same as without the disk present!

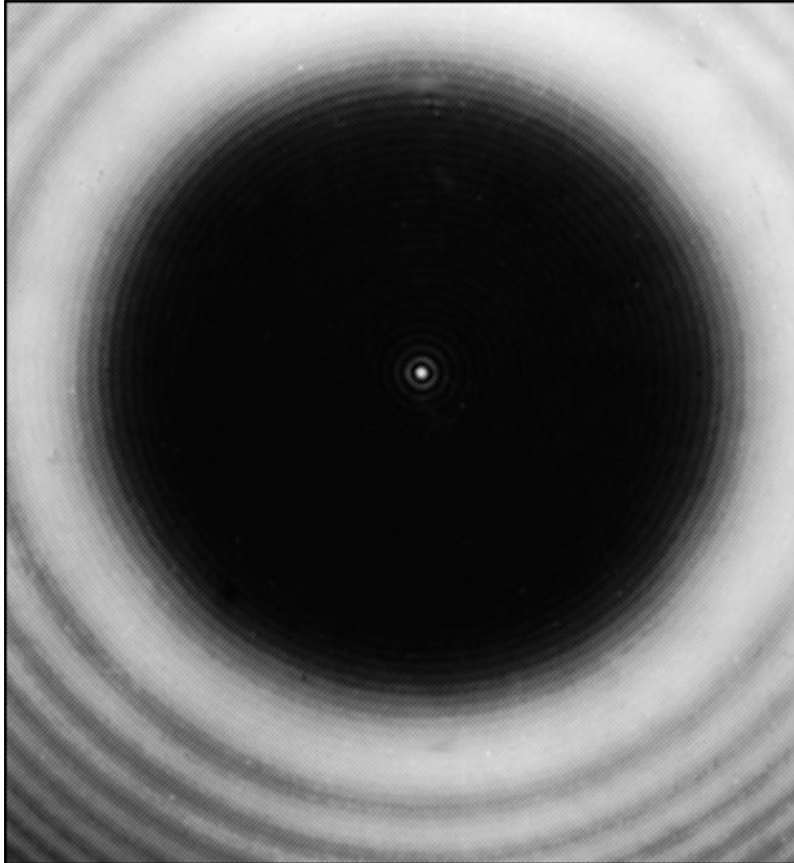


How absurd!

Siméon Denis Poisson (1781-1840)



Poisson/Arago spot

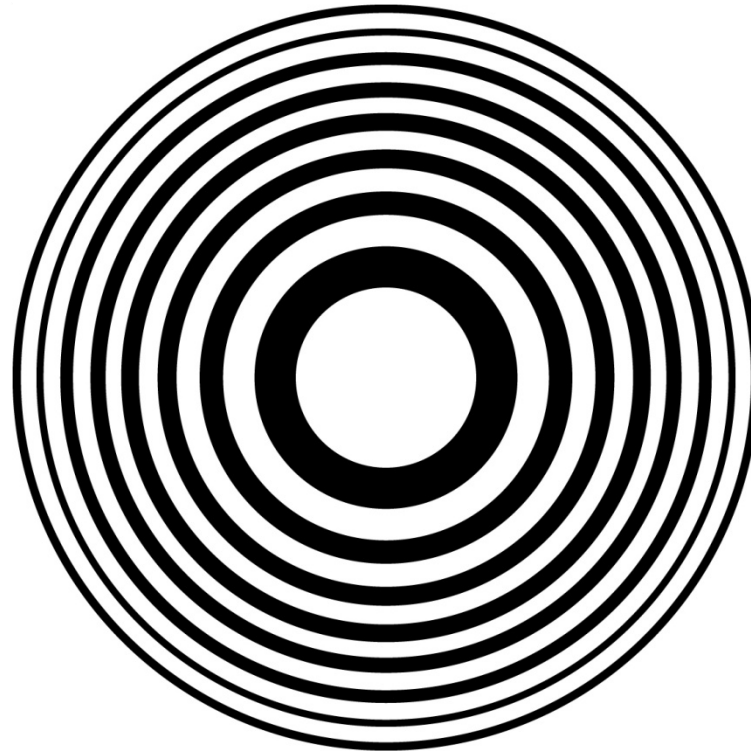


**François Arago
(1786-1853)**



The Fresnel zone plate

16 zones



$$A_n = a_1 - a_2 + a_3 - a_4 + \dots a_n$$

If the 2nd, 4th, 6th, etc. zones are blocked, then:

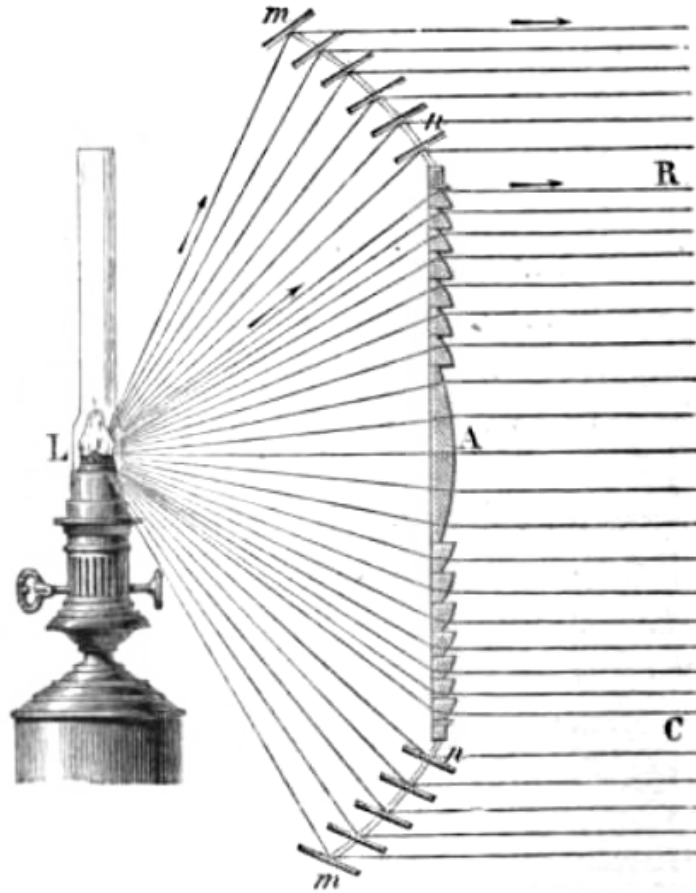
$$A_{16} = a_1 + a_3 + a_5 + a_7 + a_9 + a_{11} + a_{13} + a_{15}$$

Amplitude at P is 16 times the amplitude of $a_1/2$

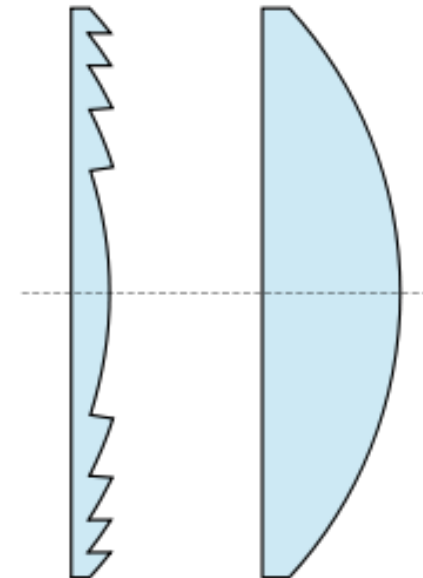
Irradiance at P is $(16)^2$ times! (a.k.a. focusing)



An alternative to blocking zones



Fresnel lens vs. *plano-convex lens*



phases of adjacent Fresnel zones changed by π



Fresnel lighthouse lens



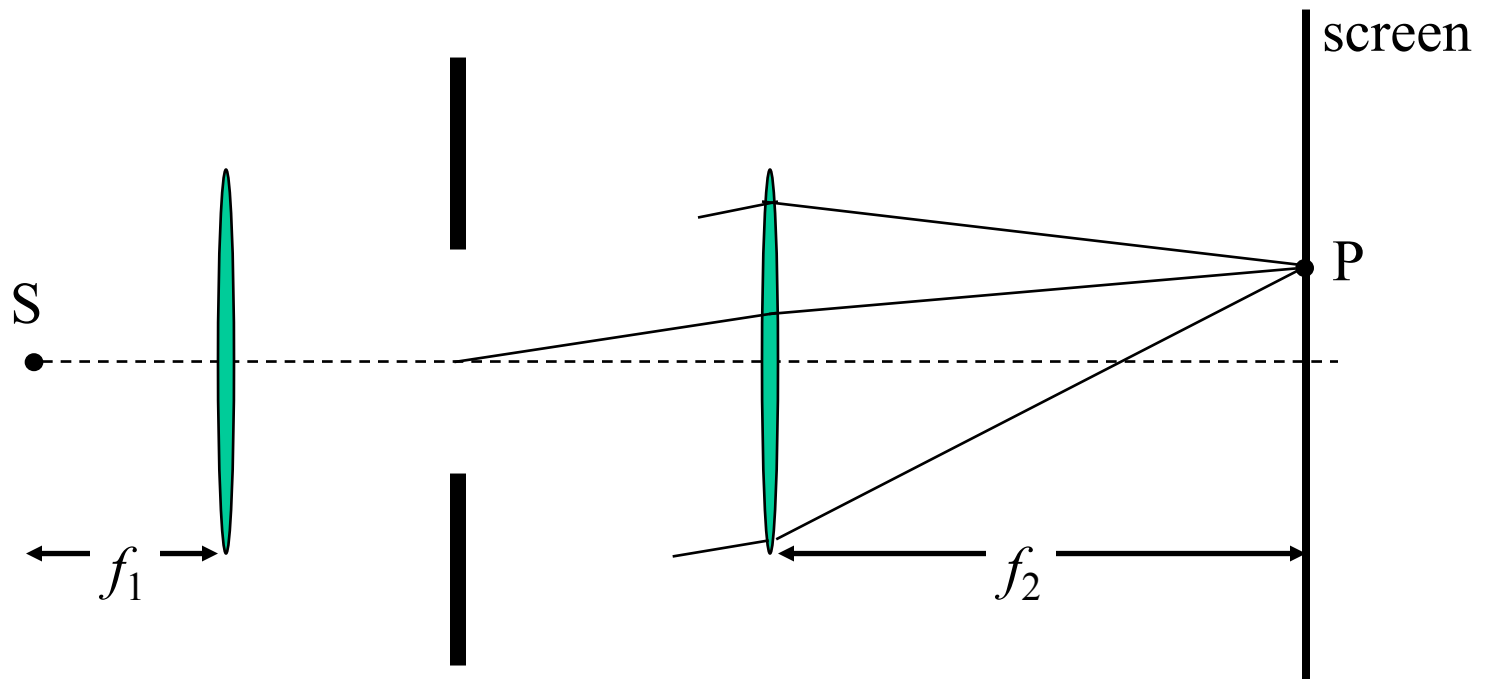
other applications:

- overhead projectors*
- automobile headlights*
- solar collectors*
- traffic lights*



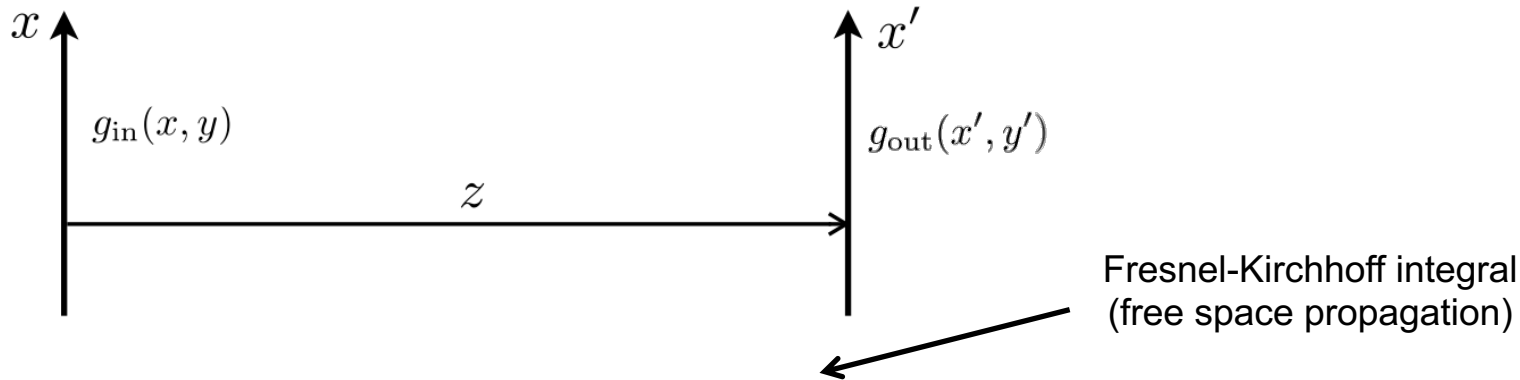
Back to Fraunhofer diffraction

- Typical arrangement (or use laser as a source of plane waves)
- Plane waves in, plane waves out





Fraunhofer diffraction



$$g_{\text{out}}(x', y'; z) = \frac{1}{i\lambda z} \exp\left\{i2\pi\frac{z}{\lambda}\right\} \iint g_{\text{in}}(x, y) \exp\left\{i\pi\frac{(x' - x)^2 + (y' - y)^2}{\lambda z}\right\} dx dy$$

$$g_{\text{out}}(x', y'; z) = \frac{1}{i\lambda z} \exp\left\{i2\pi\frac{z}{\lambda}\right\} \iint g_{\text{in}}(x, y) \exp\left\{i\pi\frac{x'^2 + x^2 - 2xx' + y'^2 + y^2 - 2yy'}{\lambda z}\right\} dx dy$$

$$\approx \underbrace{\exp\left\{i2\pi\frac{z}{\lambda} + i\pi\frac{x'^2 + y'^2}{\lambda z}\right\}}_{|\dots| = 1} \iint g_{\text{in}}(x, y) \exp\left\{-i2\pi\frac{xx' + yy'}{\lambda z}\right\} dx dy$$

$$u \equiv \frac{x'}{\lambda z} \quad v \equiv \frac{y'}{\lambda z}$$

$$g_{\text{out}}(x', y'; z) \approx \exp\left\{i2\pi\frac{z}{\lambda} + i\pi\frac{x'^2 + y'^2}{\lambda z}\right\} \iint g_{\text{in}}(x, y) \exp\left\{-i2\pi(ux + vy)\right\} dx dy$$



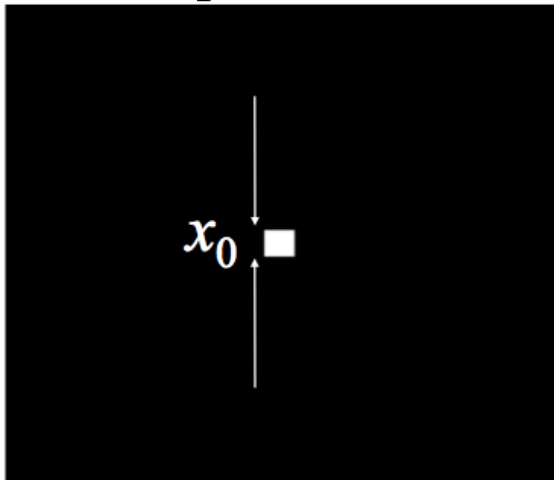
Fraunhofer diffraction \propto Fourier Transform: Rectangular aperture

Fourier transform

$$g_{\text{in}}(x, y) = \text{rect}\left(\frac{x}{x_0}\right) \text{rect}\left(\frac{y}{y_0}\right)$$
$$G_{\text{in}}(u, v) = x_0 y_0 \text{sinc}(x_0 u) \text{sinc}(y_0 v)$$

$$g_{\text{out}}(x', y'; z \rightarrow \infty) \propto \text{sinc}\left(\frac{x_0 x'}{\lambda z}\right) \text{sinc}\left(\frac{y_0 y'}{\lambda z}\right)$$

Input field

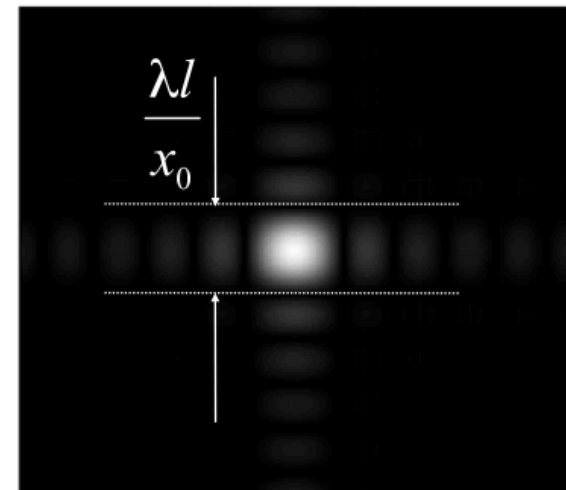


free space
propagation by

$l \rightarrow \infty$



Far-field





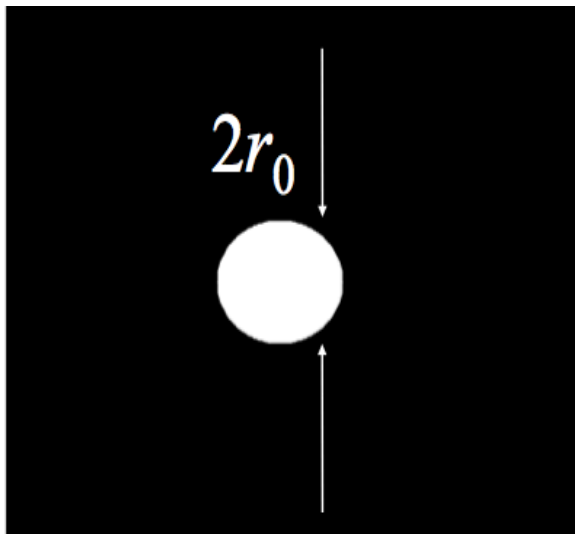
Circular aperture

$$g_{\text{in}}(x, y) = \text{circ} \left(\frac{\sqrt{x^2 + y^2}}{r_0} \right)$$

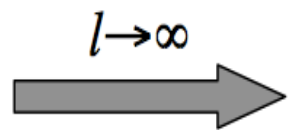
$$G_{\text{in}}(u, v) = r_0^2 \text{jinc} \left(r_0 \sqrt{u^2 + v^2} \right)$$

$$\equiv r_0 \frac{J_1 \left(2\pi \sqrt{u^2 + v^2} \right)}{\sqrt{u^2 + v^2}} \quad g_{\text{out}}(x', y'; z \rightarrow \infty) \propto \text{jinc} \left(\frac{2\pi r_0 \sqrt{x'^2 + y'^2}}{\lambda z} \right)$$

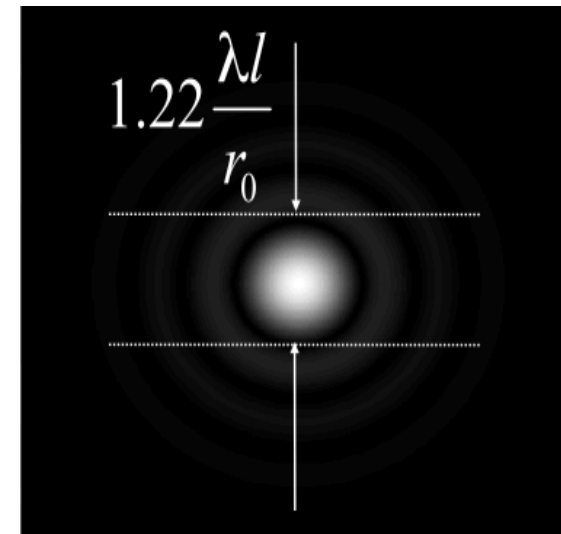
Input field



free space
propagation by



Far-field
Airy pattern





Fourier transform pair

$$G(\nu) = \int_{-\infty}^{+\infty} g(t) \exp \{ -i2\pi\nu t \} dt.$$

1D

$$g(t) = \int_{-\infty}^{+\infty} G(\nu) \exp \{ i2\pi\nu t \} d\nu.$$

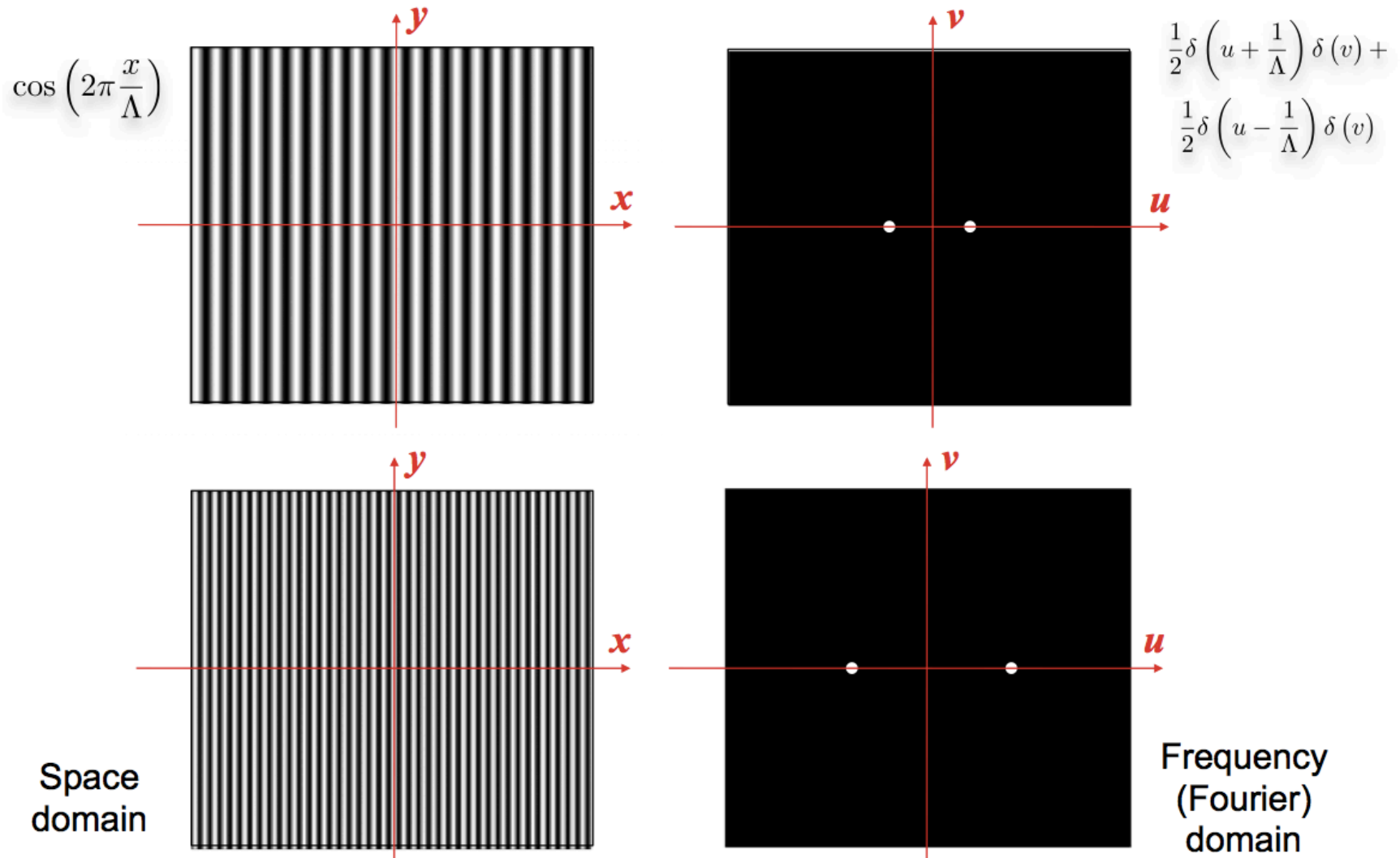
$$G(u, v) = \iint_{-\infty}^{+\infty} g(x, y) \exp \{ -i2\pi(ux + vy) \} dx dy.$$

2D

$$g(x, y) = \iint_{-\infty}^{+\infty} G(u, v) \exp \{ i2\pi(ux + vy) \} du dv.$$

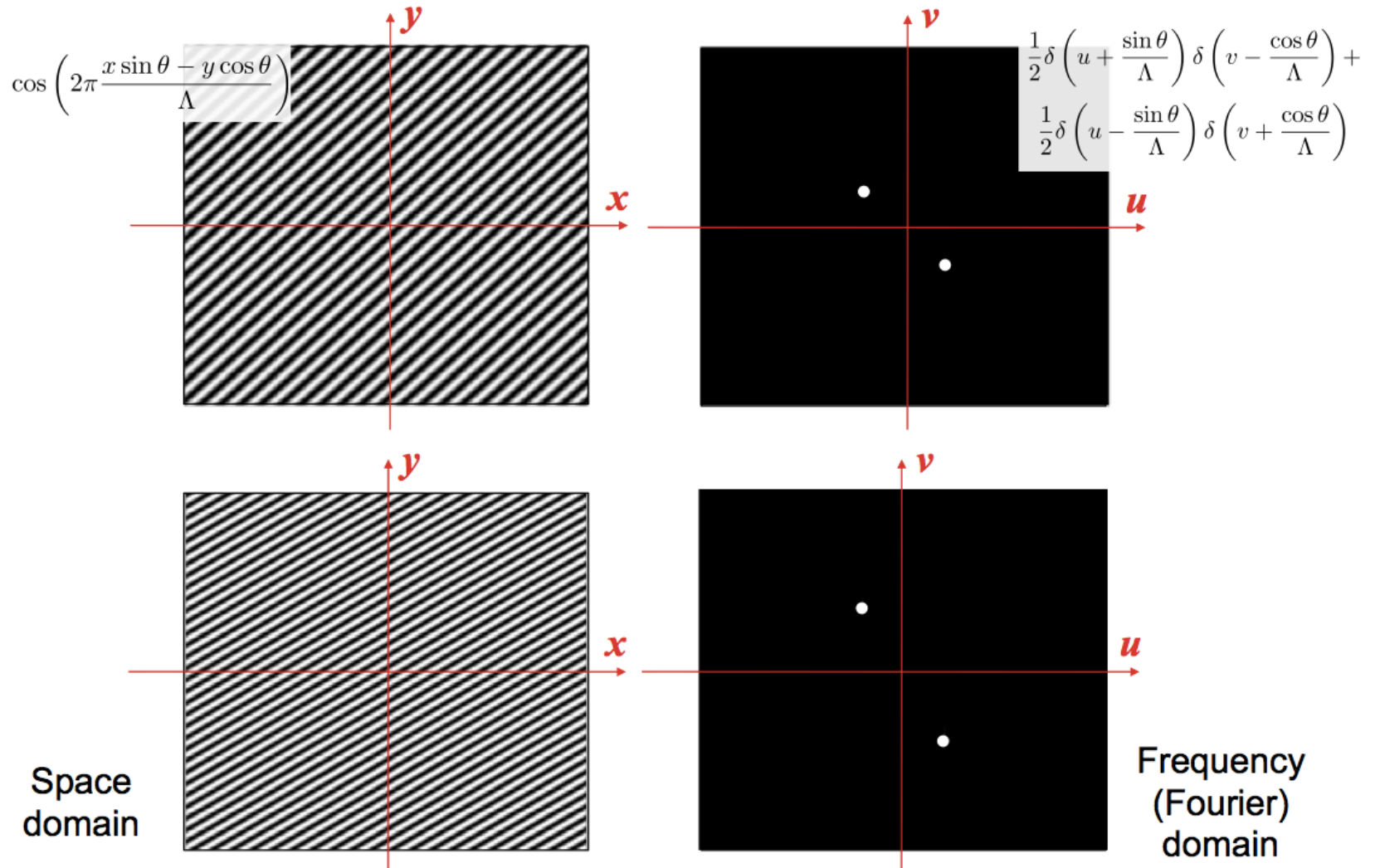


Spatial domain \leftrightarrow (angular) frequency domain





Tilted grating

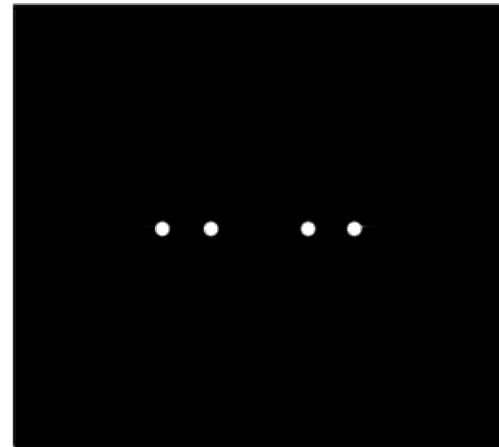
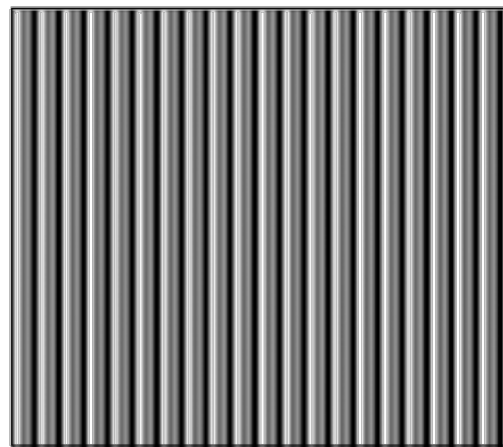
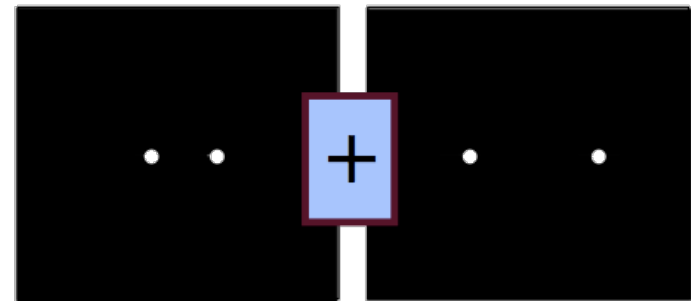
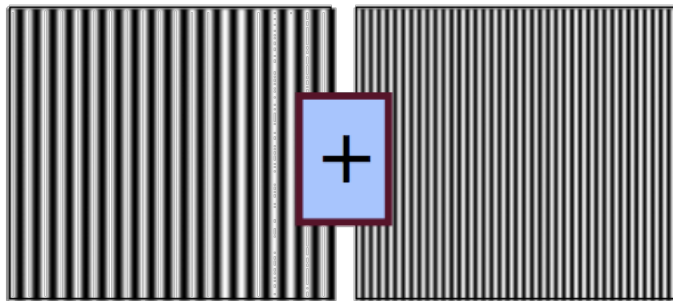




Linear superposition

$$a_1 \cos\left(2\pi \frac{x}{\Lambda_1}\right) + a_2 \cos\left(2\pi \frac{x}{\Lambda_2}\right)$$

$$\frac{a_1}{2} \delta\left(u + \frac{1}{\Lambda_1}\right) \delta(v) + \frac{a_2}{2} \delta\left(u + \frac{1}{\Lambda_2}\right) \delta(v) +$$
$$\frac{a_1}{2} \delta\left(u - \frac{1}{\Lambda_1}\right) \delta(v) + \frac{a_2}{2} \delta\left(u - \frac{1}{\Lambda_2}\right) \delta(v)$$

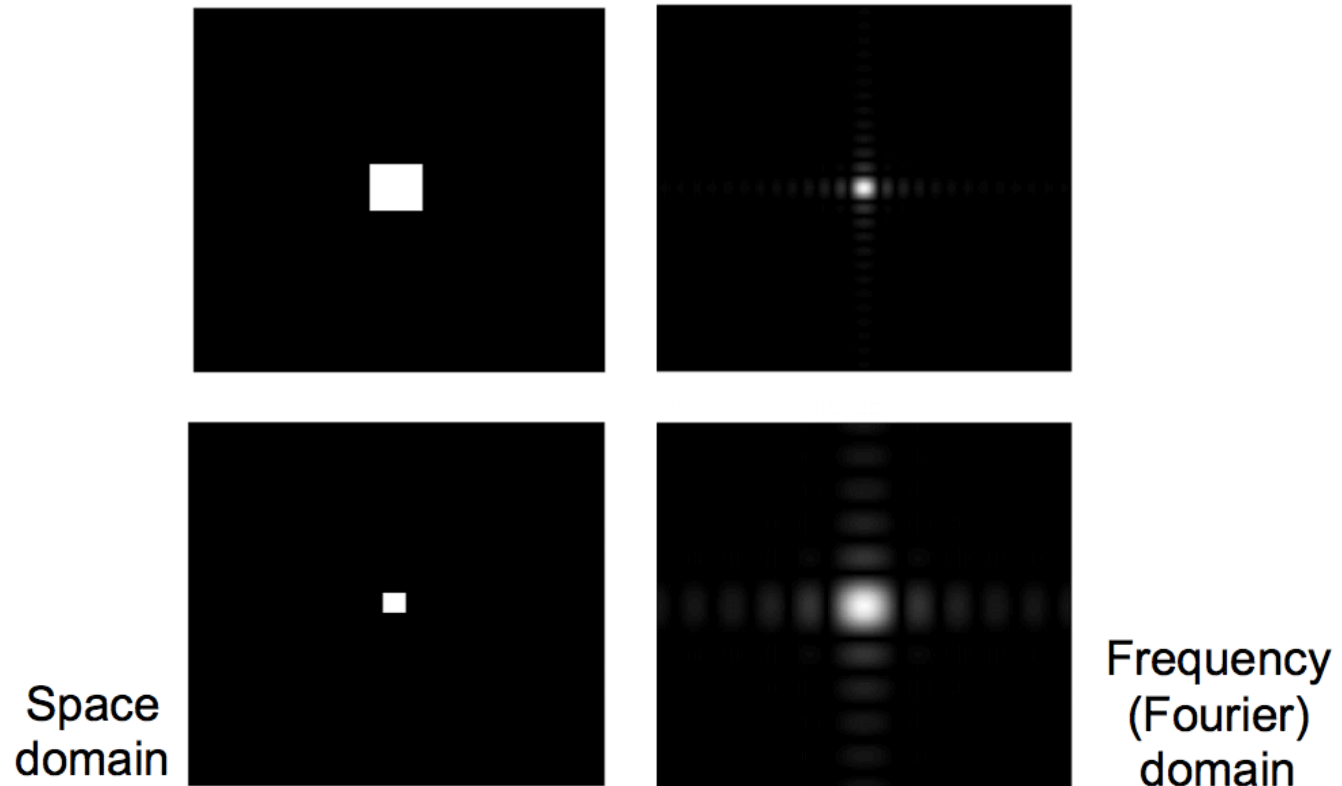


Space domain

Frequency (Fourier) domain



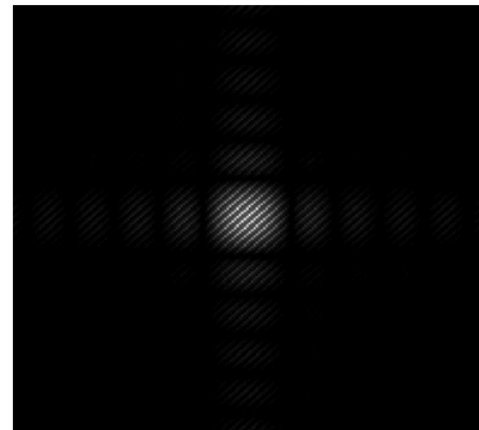
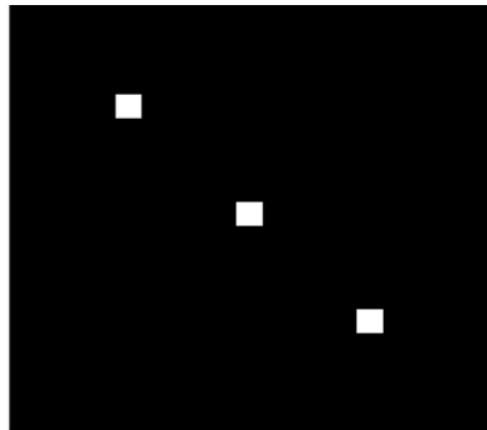
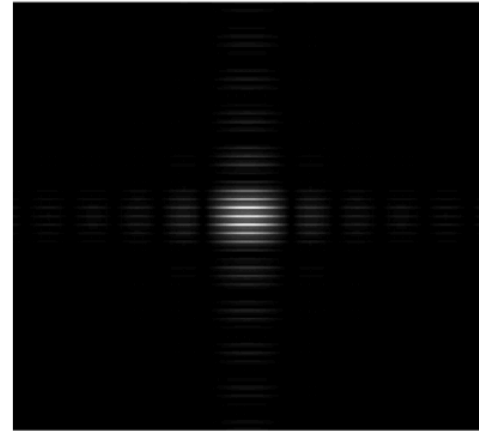
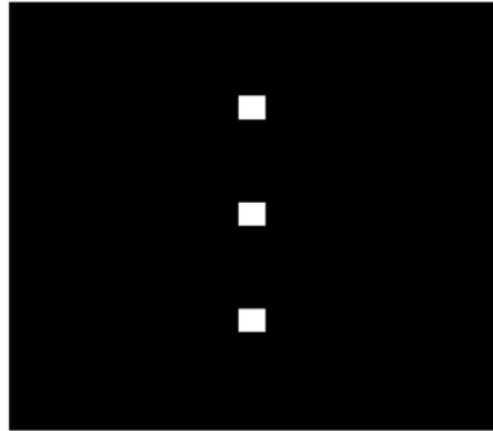
Scaling



$$\mathcal{F} \left\{ g \left(\frac{x}{a}, \frac{y}{b} \right) \right\} = |ab| G(au, bv)$$



Shift theorem



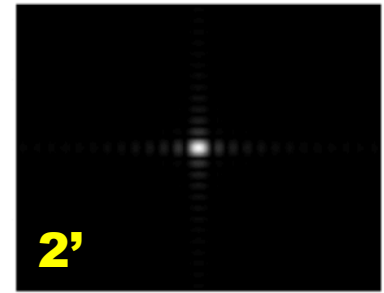
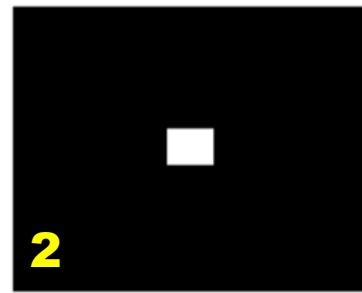
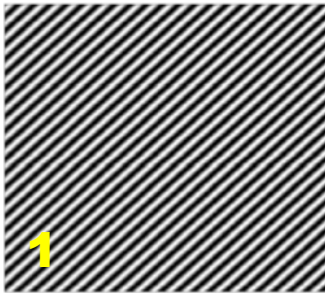
Space
domain

Frequency
(Fourier)
domain

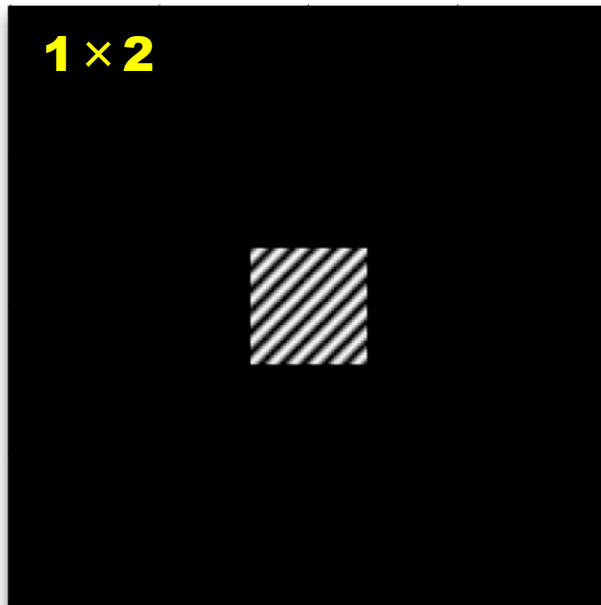
$$\mathcal{F} \{g(x - a, y - b)\} = \exp \{i2\pi(au + bv)\} G(u, v)$$



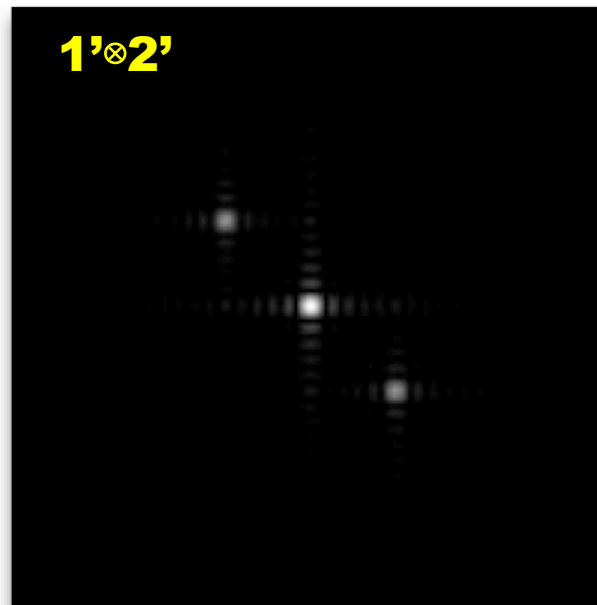
The convolution theorem



multiplication



convolution



$$\mathcal{F}\{f * g\} = \mathcal{F}\{f\} \cdot \mathcal{F}\{g\} \quad \text{or} \quad \mathcal{F}\{f \cdot g\} = \mathcal{F}\{f\} * \mathcal{F}\{g\}$$



Links/references

http://edu.tnw.utwente.nl/inlopt/overhead_sheets/Herek2010/week7/13.Fresnel%20diffraction.ppt

http://ocw.mit.edu/courses/mechanical-engineering/2-71-optics-spring-2009/video-lectures/lecture-17-fraunhofer-diffraction-fourier-transforms-and-theorems/MIT2_71S09_lec17.pdf