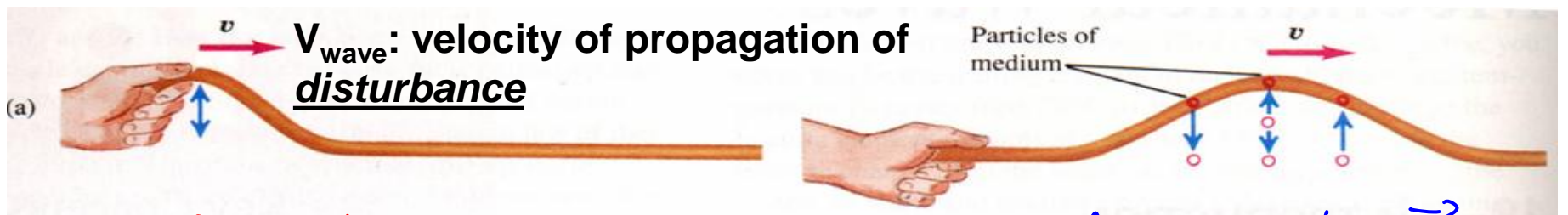


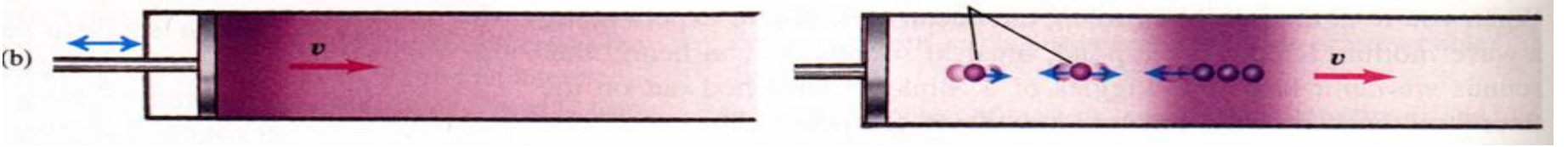
Recap: Waves

- Disturbance that propagates at wave speed v_{wave}
- Propagate energy and momentum, but not mass
- Reflected by boundaries

$\updownarrow \rightarrow \vec{v}_{wave}$ • Transverse Wave: displacement of medium \perp to \vec{v}_{wave}



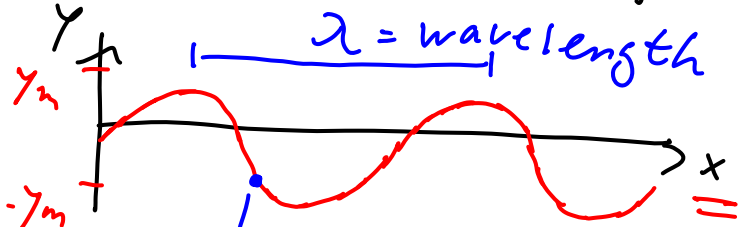
$\leftrightarrow \rightarrow \vec{v}_{wave}$ • Longitudinal Wave: displacement of medium along \vec{v}_{wave}



• Traveling sinusoidal transverse wave:

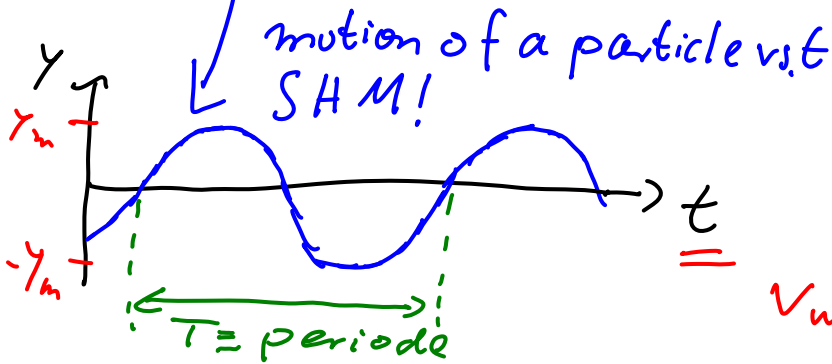
Crest "moves" sine-function motion of particles

sign gives direction of wave!
↓



$$y(x, t) = y_m \sin[kx \mp \omega t]$$

space and time dependency in argument of one sine-function!



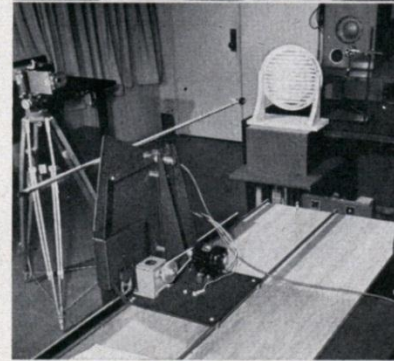
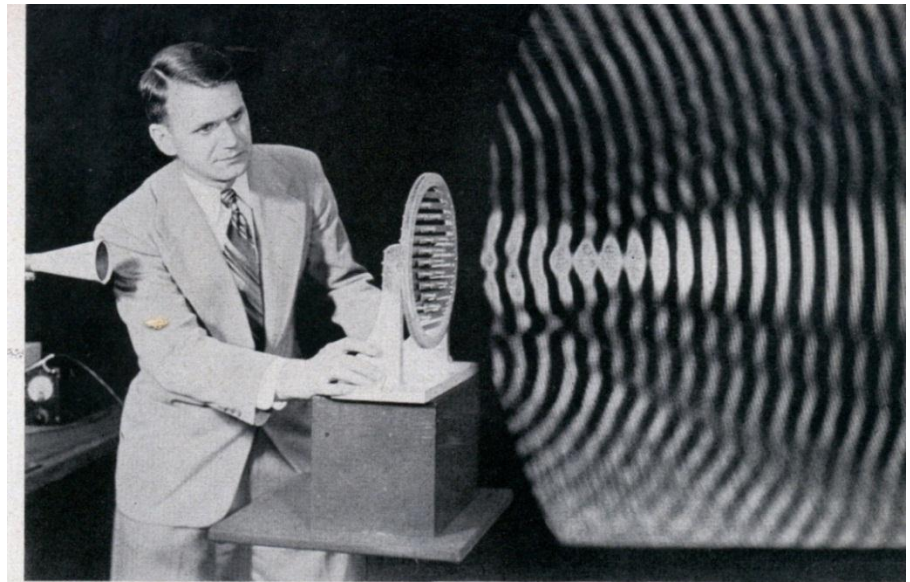
$$k = \frac{2\pi}{\lambda} = \text{wave number}$$

$$\omega = 2\pi f = \frac{2\pi}{T} = \text{angular freq.}$$

$$v_{\text{wave}} = \frac{\omega}{k} = \frac{\lambda}{T} = \lambda f = \text{"speed of crest"}$$

Today:

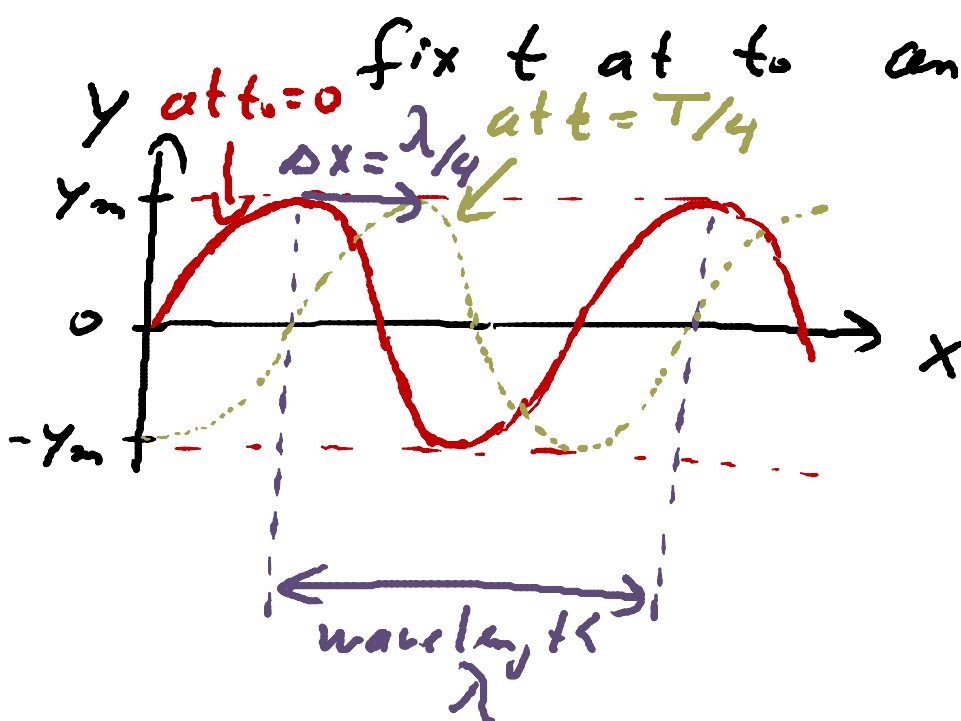
- Wave velocity
- Longitudinal traveling waves
 - Sound waves and intensity
 - Beats



Neon Lamp Traces Sound Wave's Picture

THAT'S a sound wave you see in the picture above. Here demonstrating how an acoustic lens focuses sound from a horn, the wave was made visible with the device at left—an aluminum rod with a microphone and a neon lamp at the end. A small motor swings the rod in a wide arc, scanning the area. The microphone picks up the sound and turns it into electric current to feed the lamp. Wherever the sound is strongest, the light is brightest, and the wave is traced out. A complete sound photo, such as this from Bell Labs, takes 10 minutes exposure.

→ "Snapshot" of the wave:



"-" => moves
in +x direction

$$y(x, t_0) = y_m \sin[kx - \omega t_0]$$

$$= y_m \sin[kx + \phi_0]$$

$$\text{with } \phi_0 = -\omega t_0 = -\frac{2\pi}{T} t_0$$

→ for $\Delta t = T/4$: crest moves by $\Delta x = \lambda/4$ in +x
direction

$$\Rightarrow \text{wave speed} = \text{"speed of crest"} = \frac{\Delta x}{\Delta t} = \frac{\lambda}{T} = \lambda f = \frac{\omega}{k}$$

indep. of time!

A transverse wave on a string traveling in the **+x** direction is described by

$$y(x,t) = y_m \sin [(2\pi(x/\lambda - t/T)].$$

If a **wave crest** (i.e., $y = y_m$) is located at **x=0** at some time **t**, at what **time t'** will it have **moved to x=λ**?

$$\Delta x = \lambda$$

$$\Rightarrow \Delta t = ?$$

\Rightarrow need wave speed:

$$v_{\text{wave}} = \frac{\Delta x}{\Delta t} = \frac{\lambda}{T}$$

\Rightarrow for crest to "move" $\Delta x = \lambda$ takes $\Delta t = T = \text{period}$

$$\Rightarrow t' = t + \Delta t = t + T = t + \frac{\lambda}{v_{\text{wave}}}$$

- A. $t' = t + T$
- B. $t' = t - T$
- C. $t' = t + \lambda / v$
- D. $t' = t - \lambda / v$

A transverse wave on a string traveling in the $+x$ direction is described by

$$y(x,t) = y_m \sin \left[\underbrace{(2\pi x/1\text{m} - 4\pi t/1\text{s})}_{2\pi \left(\frac{x}{1\text{m}} - \frac{t}{0.5\text{s}} \right)} \right] = y_m \sin \left[2\pi \left(\frac{x}{\lambda} - \frac{t}{T} \right) \right]$$

where x is in m and t is in s. $2\pi \left(\frac{x}{1\text{m}} - \frac{t}{0.5\text{s}} \right)$

(1) What is the **wavelength** of the wave?

$$\lambda = 1\text{m}$$

(2) What is the **period** of the local oscillations?

$$T = 1/2\text{ s}$$

(3) What is the **wave speed**?

$$v = \lambda / T = \frac{1\text{m}}{0.5\text{s}} = 2\text{ m/s}$$

The **speed v** of a wave on a **stretched string** or wire depends on the **wire tension F_T** and the **wire's mass per unit length μ** .

From **dimensional analysis**, what must be the **relation** between v , F_T , and μ ?

$$\begin{array}{l} v_{\text{wave}} \\ \frac{m}{s} \end{array} \quad \left| \quad \begin{array}{l} F_T \\ N = \frac{kg \cdot m}{s^2} \end{array} \quad \mu = \frac{\text{mass}}{\text{length}} \\ \quad \frac{kg}{m}$$

$$\Rightarrow v_{\text{wave}} \propto \sqrt{\frac{F_T}{\mu}}$$

- A. $v \propto F_T / \mu$
- B. $v \propto (F_T / \mu)^{1/2}$**
- C. $v \propto \mu / F_T$
- D. $v \propto (\mu / F_T)^{1/2}$

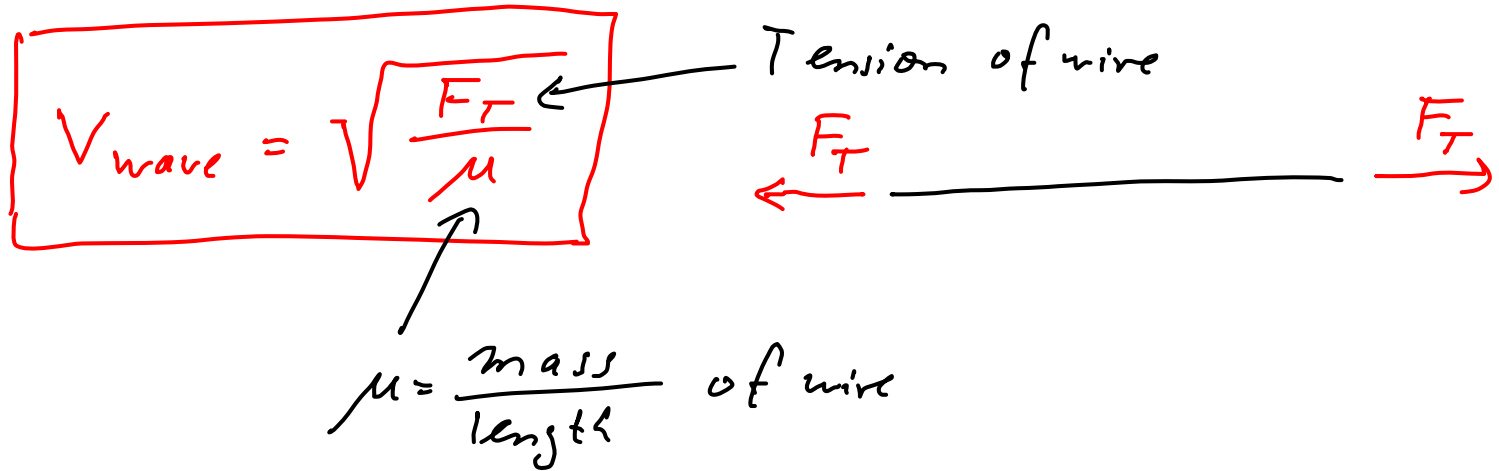
→ wave speed on a stretched string / wire:

$$V_{\text{wave}} = \sqrt{\frac{F_T}{\mu}}$$

Tension of wire

F_T ← ————— → F_T

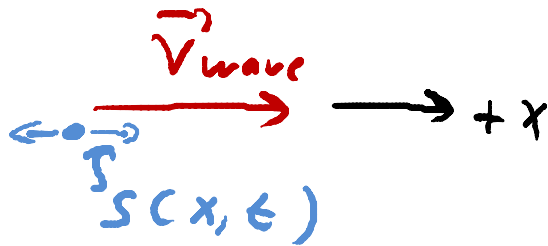
$\mu = \frac{\text{mass}}{\text{length}}$ of wire

A diagram illustrating the wave speed on a stretched string. On the left, a red-bordered box contains the equation $V_{\text{wave}} = \sqrt{\frac{F_T}{\mu}}$. An arrow points from the text 'Tension of wire' to the F_T term in the numerator of the equation. Below the box, an arrow points from the text ' $\mu = \frac{\text{mass}}{\text{length}}$ of wire' to the μ term in the denominator. To the right of the box, a horizontal line represents a stretched wire. Two red arrows labeled F_T point outwards from the ends of the wire, representing the tension forces.

⇒ change tension to tune violin via changing V_{wave} !

→ Sound Waves (traveling)

- Longitudinal waves (in gas or liquid)



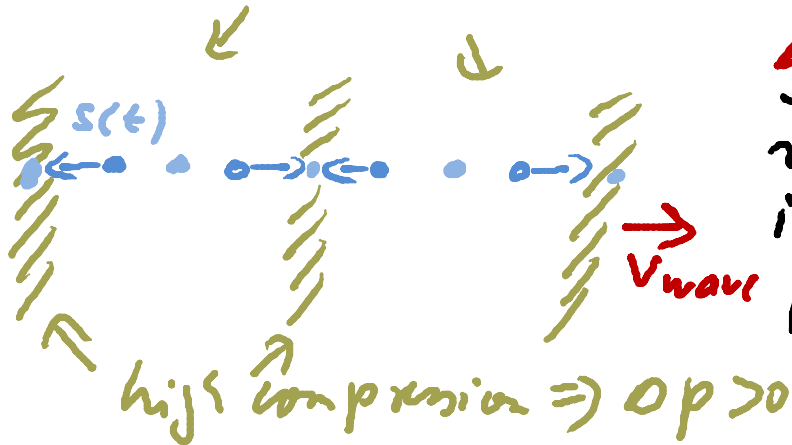
Displacement of molecules back and forth along x (\parallel to \vec{V}_{wave})

$$s(x, t) = s_m \cos [kx \mp \omega t] \left. \begin{array}{l} \text{dipl.} \\ \text{wave} \end{array} \right\}$$

↑
displacement of molecule at position x and time t

⇒ Displacement wave causes compression and expansion of gas / liquid ⇒ pressure wave

expansion: $\Delta p < 0$



$$\Delta p(x, t) = \Delta p_m \sin [kx \mp \omega t]$$

↑
relative to p_0 ,
i.e. change in
pressure due
to wave

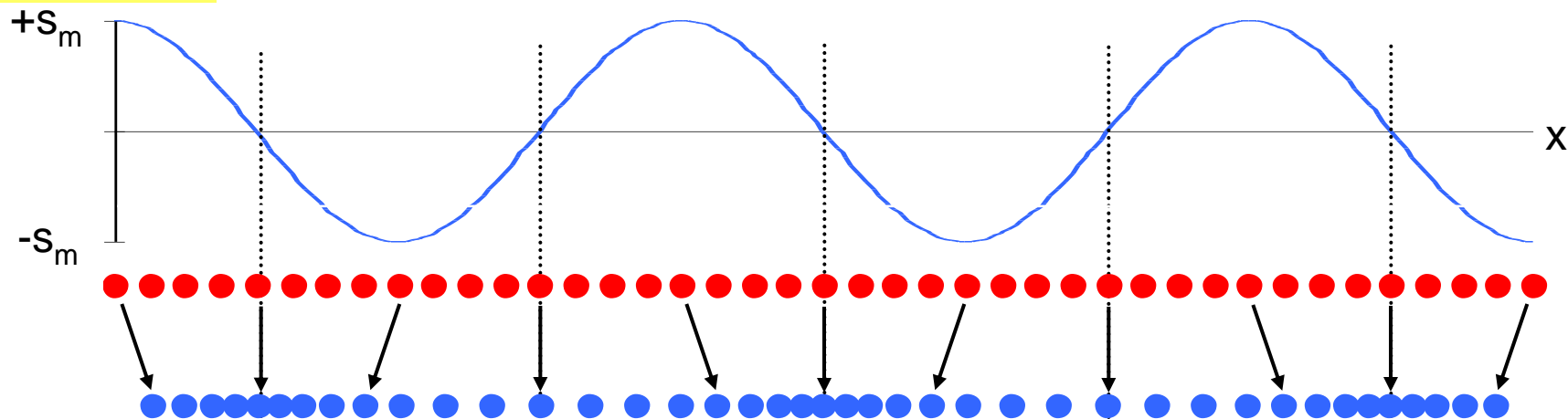
↑
90° out of phase
with displacement
wave!

Sound Waves:

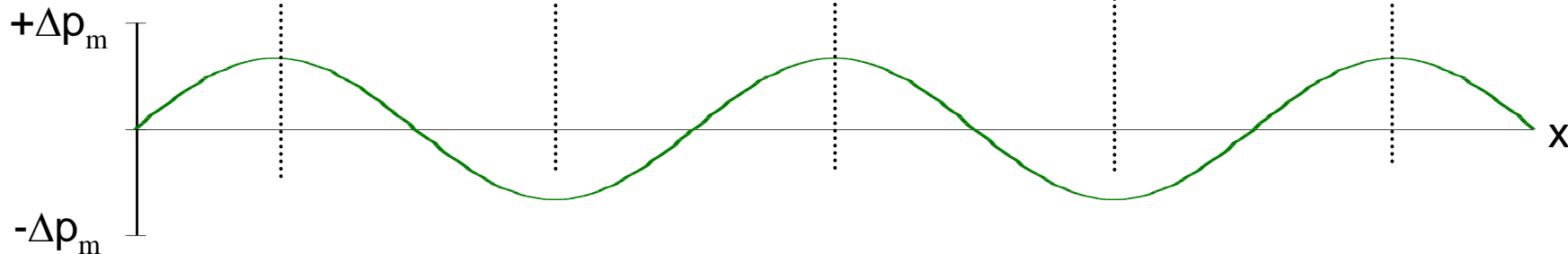
longitudinal molecular displacement: $s(x,t) = s_m \cos[kx - \omega t]$

excess pressure: $\Delta p(x,t) = \Delta p_m \sin [kx - \omega t]$

$s(x,t=0):$

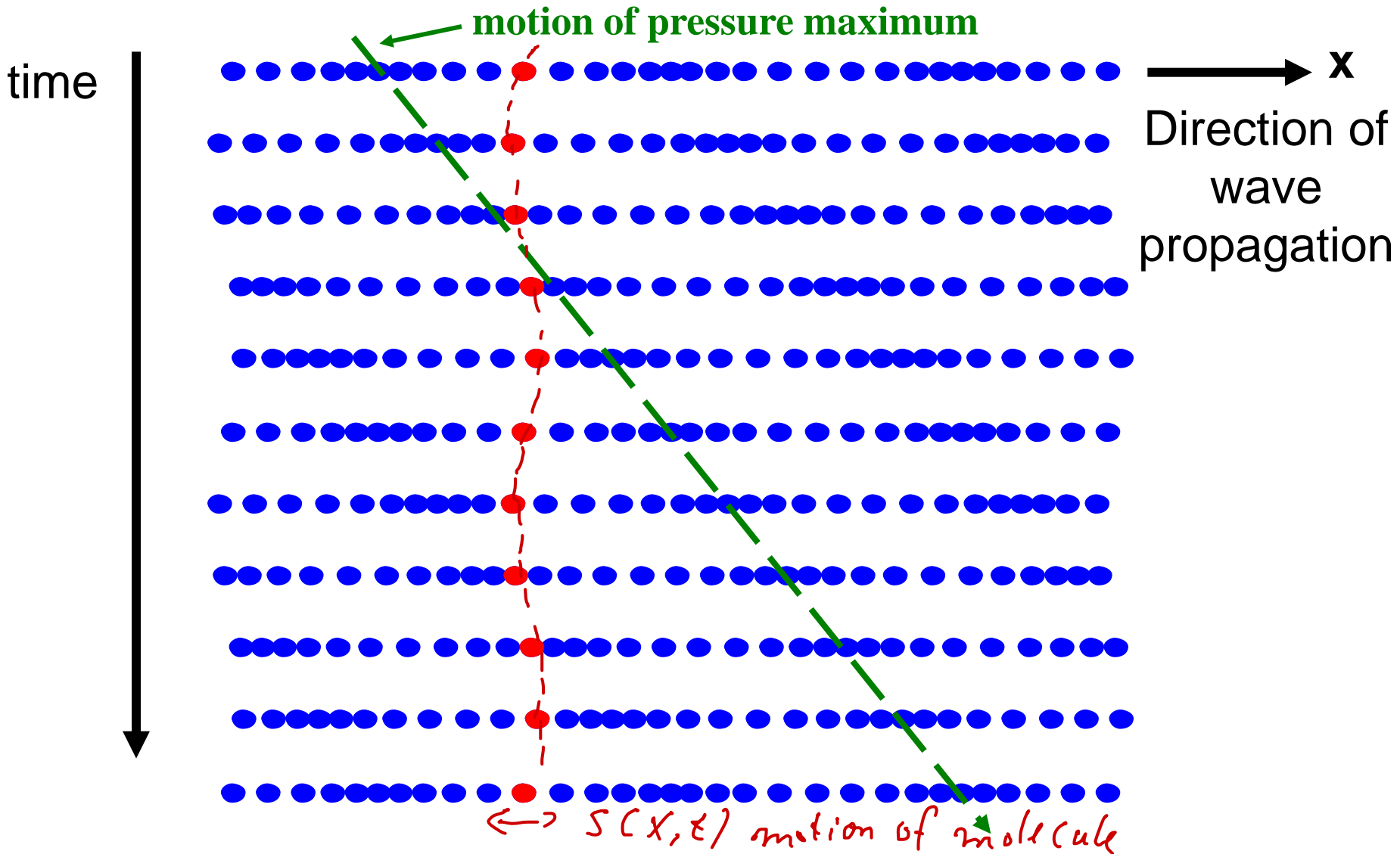


$\Delta p(x,t=0):$



Sound Waves:

longitudinal molecular displacement: $s(x,t) = s_m \cos[kx - \omega t]$



⇒ pressure amplitude ΔP_m and displacement amplitude S_m are related:

$$\Delta P_m = (V_{\text{sound wave}} \rho \omega) S_m$$

↑ density of medium

⇒ Speed of sound wave:

$$V_{\text{sound wave}} = \sqrt{\frac{B}{\rho}}$$

B: bulk modulus:

$$\frac{F}{A} = -B \frac{\Delta V}{V}$$

↑

Examples:

$$\begin{aligned} V_{\text{sound}} &= 343 \text{ m/s in air } (\approx \frac{1}{5} \text{ mile/s}) \\ &= 1482 \text{ m/s in water} \\ &= 6420 \text{ m/s in aluminium} \end{aligned}$$

→ Intensity I of a sound wave:

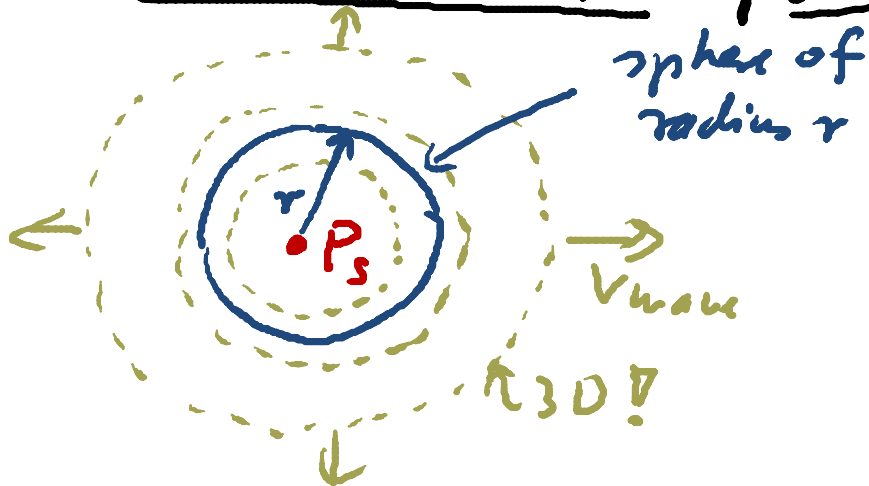
$I = \frac{\text{energy transferred by a sound wave per unit time}}{\text{area through which it is transferred}}$

= $\frac{\text{acoustic power}}{\text{area}} = \frac{P}{A}$ $[I] = \frac{W}{m^2}$

$I = \left(\frac{1}{2} \rho \omega^2 S_m^2 \right) v_{\text{wave}} = \frac{\Delta P_m^2}{2 v_{\text{wave}} \rho}$

\uparrow density of medium $\underbrace{\quad}_{v_{\text{max}}^2}$

→ for an isotropic point source:



$I = \frac{P_s}{A} = \frac{P_s}{4\pi r^2}$ } area of sphere

A **point source** emits sound waves isotropically. The **intensity** of the wave **1 m** from the source is **1 W/m²**.

What is the **intensity** of the wave **2 m** from the source?

$$I = \frac{P_s}{4\pi r^2} \propto \frac{1}{r^2}$$

$$\Rightarrow \frac{I(2\text{m})}{I(1\text{m})} = \left(\frac{1\text{m}}{2\text{m}}\right)^2 = \frac{1}{4}$$

$$I(2\text{m}) = ?$$

- A. 0.25 W/m²**
- B. 0.5 W/m²**
- C. 1 W/m²**
- D. 2 W/m²**
- E. 4 W/m²**

→ Sound level β :

range of human hearing : $I = 10^{-12} \frac{W}{m^2} \dots 1 \frac{W}{m^2}$
threshold of hearing threshold of pain

⇒ use logarithmic scale : decibel scale

$$\beta = (10 \text{ dB}) \cdot \log_{10} \left(\frac{I}{I_0} \right)$$

$$I_0 = 10^{-12} \frac{W}{m^2} \text{ (reference intensity)}$$

"units" = decibel (dB) = dimensionless

"dB" means that we have taken $10 \cdot \log_{10}(\text{quantity})$

$$\begin{array}{l} I = 10^{-12} \frac{W}{m^2} \rightarrow \beta = 0 \text{ dB} \\ I = 1 \frac{W}{m^2} \rightarrow \beta = 120 \text{ dB} \end{array} \left. \vphantom{\begin{array}{l} I = 10^{-12} \frac{W}{m^2} \\ I = 1 \frac{W}{m^2} \end{array}} \right\} \begin{array}{l} \text{factor of 10 in} \\ \text{intensity every} \\ + 10 \text{ dB} \end{array}$$

→ Beats:

Two sound waves of equal amplitude but different frequencies, detected at some position $x=0$

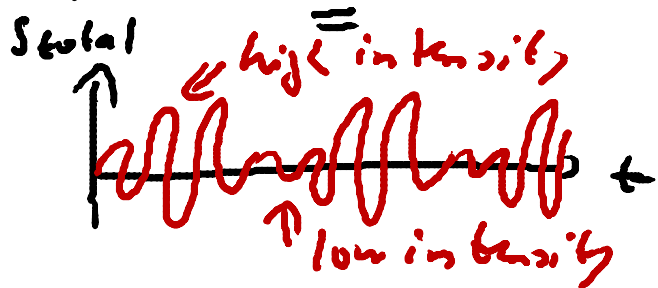
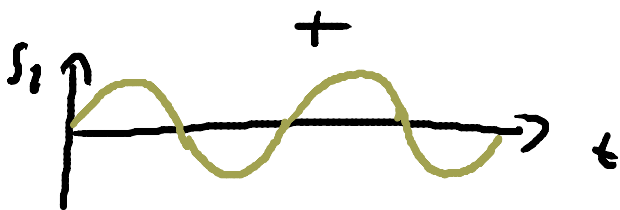
$$S_1 = \dots$$

$$S_2 = \dots$$

$$S_1(t) = S_m \cos(\omega_1 t)$$

$$S_2(t) = S_m \cos(\omega_2 t)$$

displacements due to waves at position $x=0$



⇒ total displacement:

(principle of superposition)

$$S_{\text{total}}(t) = S_1(t) + S_2(t)$$

$$= S_m [\cos(\omega_1 t) + \cos(\omega_2 t)]$$

$$= 2 S_m [\cos(\omega' t) \cdot \cos(\omega_{\text{avg}} t)]$$

$$\text{with } \omega' = \frac{|\omega_1 - \omega_2|}{2} = \frac{\text{difference}}{2}$$

$$\omega_{\text{avg}} = \frac{\omega_1 + \omega_2}{2} = \text{average of } \omega\text{'s}$$

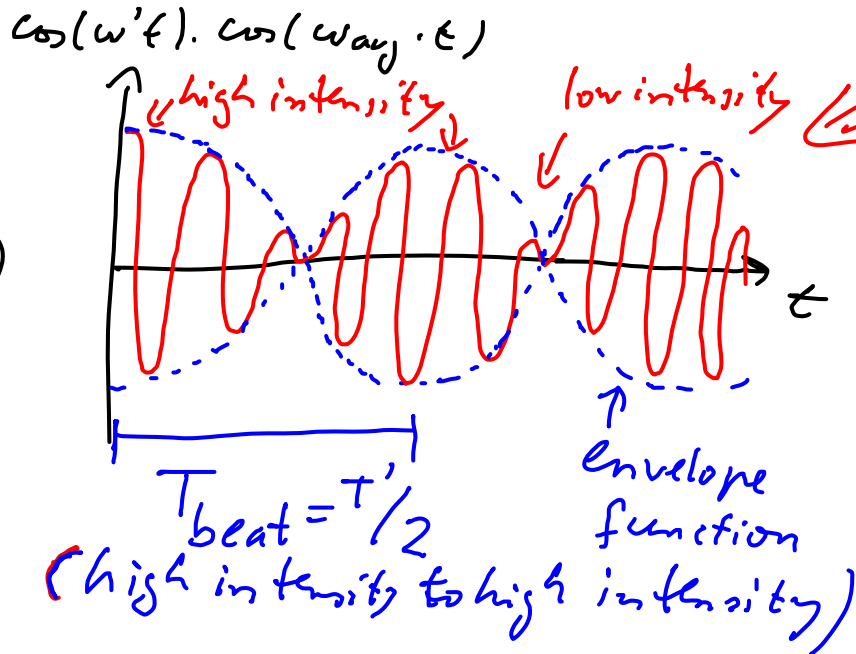
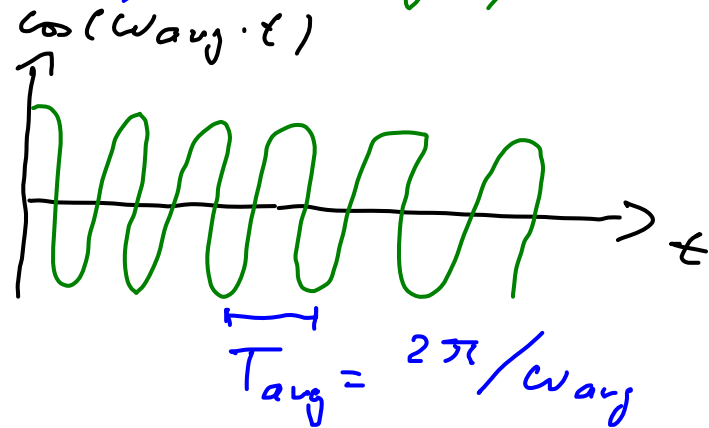
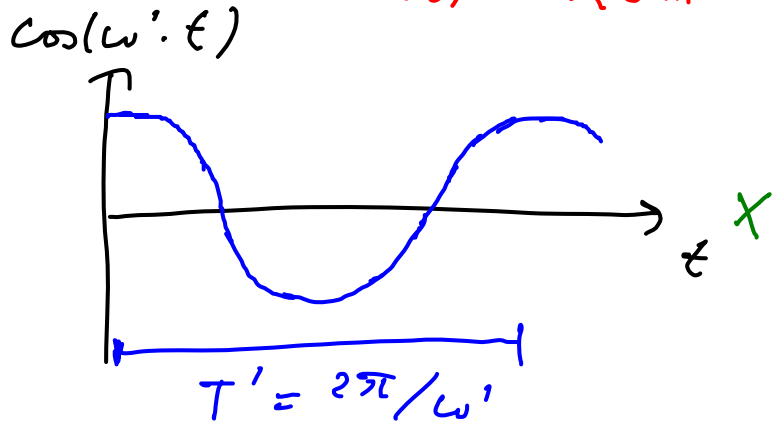
$$\cos(A) + \cos(B)$$

$$= 2 \cos\left(\frac{A+B}{2}\right) \cdot \cos\left(\frac{A-B}{2}\right)$$

\Rightarrow for ω_1 close to (but not identical to) ω_2 :

$$\omega_1 \sim \omega_2 \quad \Rightarrow \quad \omega_{avg} \gg \omega'$$

$$S(t) = 2 S_m \cos(\omega' t) \cos(\omega_{avg} t)$$



beats! Intensity oscillates in time!

$$f_{beats} = \frac{1}{T_{beats}} = \frac{1}{T'/2} = 2f' = \frac{2\omega'}{2\pi}$$

$$= \frac{|\omega_1 - \omega_2|}{2\pi} = \underline{\underline{|f_1 - f_2|}}$$

= |difference between the two frequencies|