

1-LOOP TECHNIQUES FOR SCATTERING AMP (Dean) 15 Apr.

1. a) SPINOR - HELICITY REVIEW
b) TREE-LEVEL BCFW REVIEW
2. INTEGRAL REDUCTION
3. UNITARITY METHOD

SEE ZUR BERN LECTURES
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Tree stuff - review of last fall

COLOR ORDERING :

$$A_n^{\text{tree}}(\{k_i, \lambda_i, a_i\}) = \sum_{\sigma \in S_n / \mathbb{Z}_n} \text{Tr}[T^{a_{\sigma(1)}} \dots T^{a_{\sigma(n)}}] \\ \times \underbrace{A_n^{\text{tree}}(k_{\sigma(1)}, \dots, k_{\sigma(n)})}_{\text{COLOR ORDERED}}$$

Parke-Taylor

$$A_3^{\text{tree}}(1^-, 2^-, 3^+) = \frac{i \langle 12 \rangle^3}{\langle 23 \rangle \langle 31 \rangle}$$

$$A_4^{\text{tree}}(1^-, 2^-, 3^+, 4^+) = \frac{i \langle 12 \rangle^4}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle 41 \rangle}$$

REMINDEES: spinor helicity

EMBED EVERYTHING IN $SL(2, \mathbb{C})$ RATHER THAN $SO(3, 1)$
FOR SOME k^μ w/ $k^2 = 0$, WRITE $\lambda \in SL(2, \mathbb{C})$

$$\lambda_k^a \lambda_k^{\dot{a}} = k^\mu \sigma_\mu^{a\dot{a}}$$

\uparrow RH
 \uparrow LH
 \uparrow $\lambda^{\dot{a}} = (\lambda^a)^\dagger$; λ not Grassmannian

$$k^\mu = \frac{1}{2} \lambda_k^a \sigma_\mu^{a\dot{a}} \lambda_k^{\dot{a}}$$

AN ALTERNATE, CONVENIENT NOTATION:

$$\langle ij \rangle = \lambda_i^a \lambda_{aj}$$

$$[ij] = \lambda_{i\dot{a}} \lambda_{\dot{a}j}$$

$$\langle ij \rangle^* = [ji]$$

\uparrow
USEFUL FOR SIMPLE CONTRACTIONS.

ANOTHER NOTATION FOR MORE COMPLICATED CONTR. (DIXON)

$$|i^+\rangle \equiv \lambda_{ia}$$

$$\langle i^-| \equiv \lambda_i^{\dot{a}}$$

$$|i^-\rangle \equiv \lambda_i^{\dot{a}}$$

$$\langle i^+| \equiv \lambda_{ia}$$

then: $k^\mu = \frac{1}{2} \langle k^- | \sigma^\mu | k^- \rangle$

$$\langle ij \rangle = \langle i^- | j^+ \rangle$$

Real power: can also write polarization in the same way:

$$\epsilon_{\mu}^{\pm}(k) = \frac{\pm \langle q^{\mp} | \sigma^{\mu} | k^{\mp} \rangle}{\sqrt{2} \langle q^{\mp} | k^{\pm} \rangle}$$

↑
where q is a reference momentum, $q^2 = 0$

IF YOU'RE CLEVER ENOUGH TO CHOOSE q INTELLIGENTLY (GAUGE CHOICE), THINGS SIMPLIFY.

eg. $A(1^-, 2^+, 3^+, \dots, n^+) = 0$

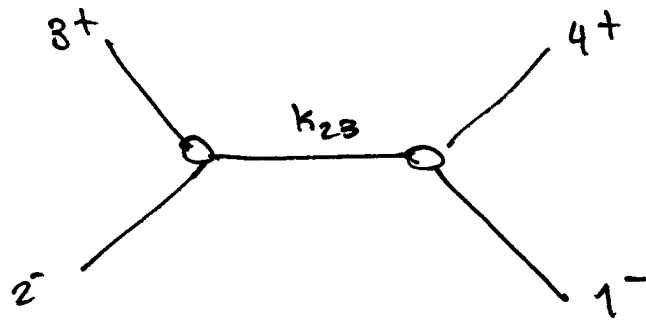
LOTS OF IDENTITIES (Shouten \leftrightarrow Jacobi)

eg. IN ANY FRAME, PICK $\epsilon \perp k$
BUT THEN BOOST FRAME & ϵ PICKS UP
A t COMPONENT. CHOICE OF q RELATED
TO THIS UNPHYSICAL PIECE.

RECALL: COLOR ORDERING (THIS PAGE)

Parke-Taylor (MHV)

$$A^{\text{tree}}(1^-, 2^-, 3^+, 4^+) =$$



$$= A_3 \frac{1}{k_{23}^2} A_3$$

Feynman: $k_{23}^2 \neq 0$ (off shell)

BCFW: SHIFT k_1, k_2 INTO \mathbb{C} PLANE,
KEEPING THEM ON-SHELL: $k_1^2 = k_2^2 = 0$

$$\begin{aligned} \text{eg } k_1^\mu &\rightarrow k_1^\mu + \frac{z}{2} \langle 1^- | \sigma^\mu | 2^- \rangle \\ k_2^\mu &\rightarrow k_2^\mu - \frac{z}{2} \langle 1^- | \sigma^\mu | 2^- \rangle \end{aligned}$$

still square to zero (can check)

WHAT DOES THIS DO? k_{23} ALSO SHIFTED.

$$0 = \oint \frac{A(z)}{z} dz \quad (\text{BODY OF } \mathbb{C})$$

$$= \boxed{A(0)} + \frac{A(\text{pole } k_{23}^2 = 0)}{z}$$

↑
what we want
ON SHELL DECOMP.

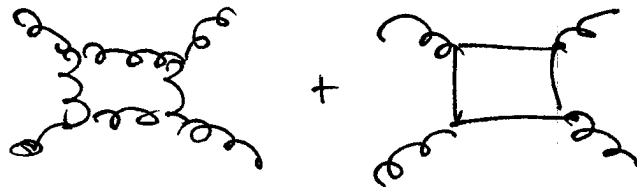
Moral: CLEVER ANALYTICITY TRICK

term by term you satisfy Ward Identity
(don't req. large cancellations like Feynman method)

LOOP LEVEL

(VERY ABBREVIATED)

Think about:



FIRST HAVE TO ORDER ORDER. [WE'LL SKIP]

ANSWER: GOES LIKE TREE LEVEL "SQUARED"

$$A_n^{1\text{-loop}} \{k_i, \lambda_i, q_i\} = \sum_c \sum_\sigma G_{n,c}(\sigma) A_{n,c,\sigma}^{1\text{-loop}}(\epsilon)$$

↑

$c=1$: tree

$c=2$: $\text{Tr}(1 \dots c) \text{Tr}(c \dots n)$

GIVES GAUGE STRUCTURE @ LOOP LEVEL

$$A_{4,1} = \text{[Self-energy diagram]} + \text{[Vertex correction diagram]} \times \frac{N_f}{N_c} \quad (\text{FACT})$$

Integral Reduction

n -point Gluon amplitude

WHICH LOOPS ARE IMPORTANT & WHICH AREN'T?

$$\begin{aligned}
 \mathcal{A}_n^{1\text{-loop}} &= \sum_{i_1 < i_2 < i_3 < i_4} d_{i_1 \dots i_4}^0 \quad \left[\text{Diagram: Box loop with external legs } i_1, i_2, i_3, i_4 \right] \\
 &+ \sum_{i_1 < i_2 < i_3} c_{i_1 \dots i_3}^0 \quad \left[\text{Diagram: Triangle loop with external legs } i_1, i_2, i_3 \right] \\
 &+ \sum_{i_1 < i_2} b_{i_1 i_2}^0 \quad \left[\text{Diagram: Bubble loop with external legs } i_1, i_2 \right] \\
 &+ \text{higher } \mathcal{O}, \text{ non div.}
 \end{aligned}$$

↑
IT'S ENCODED LOOP FUNCTIONS

b, c, d: RATIONAL FUNC OF EXTERNAL MOMENTA

(WE DON'T KNOW THEM)

↑
WANT TO FIND THEM

PREVIOUSLY: USED RES. THM TO DETERMINE UNSHIFTED
P.E FROM SHIFTED ONE

NOW: USE OPTICAL THM.

WE WILL CUT THESE & IMPOSE THAT DISCONTINUITY MATCHES; this will fix b, c, d

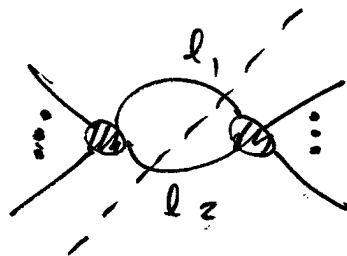
↳ we will get recursion relations, cut-by-cut

3. UNITARITY METHOD ("analyticity method")

@ the INTEGRAND LEVEL :

$$\begin{aligned}
 A_n^{1\text{-loop}}(\ell) &= \sum \frac{\bar{d}(\ell)}{D_1 D_2 D_3 D_4} \\
 &+ \sum \frac{\bar{c}(\ell)}{D_1 D_2 D_3} \\
 &+ \sum \frac{\bar{b}(\ell)}{D_1 D_2}
 \end{aligned}$$

OPTICAL THM :



CAN EXTEND THE OPTICAL THM TO n-CUTS :

$$\text{Disc}_{s_{i_1} \dots s_{i_k}} A_n^{1\text{-loop}} = \sum_{\substack{\text{States} \\ \uparrow \\ \text{ON SHELL} \\ \text{INTERVAL}}} d\Phi \underset{\substack{\uparrow \\ \text{PHASE} \\ \text{SPACE}}}{A_{n_1}^{\text{tree}} \dots A_{n_k}^{\text{tree}}}$$

end up w/ recursion relations; eg.

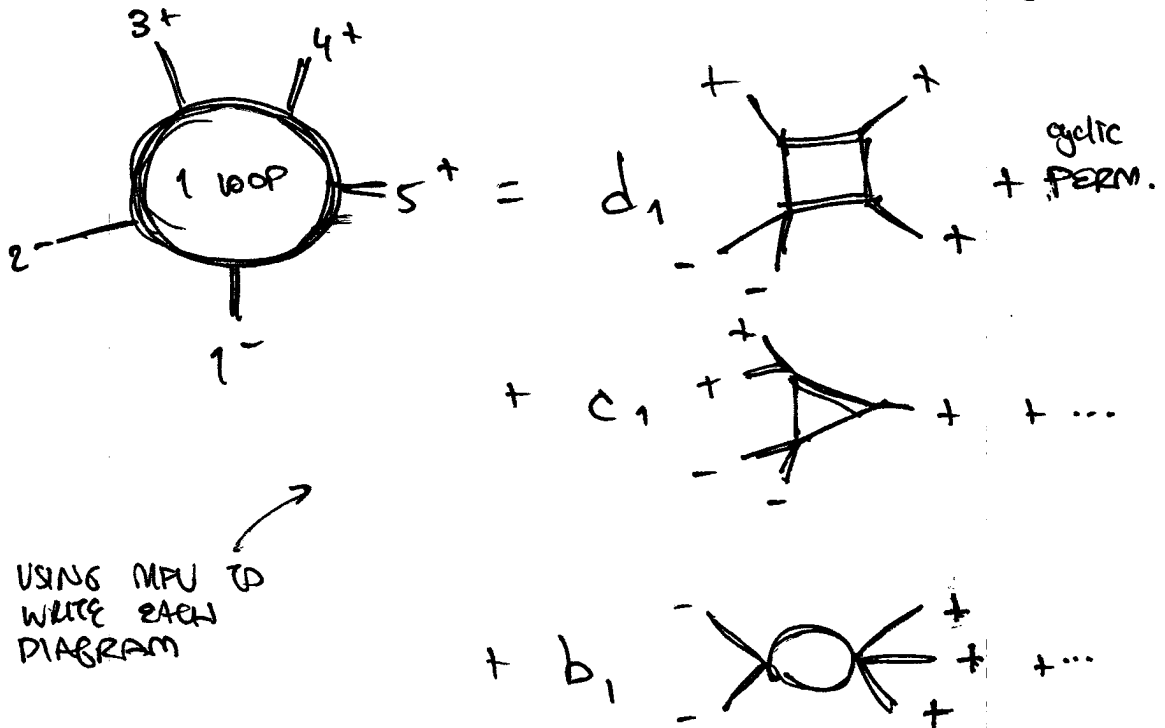
$$\bar{J}_{i_1 \dots i_4}(\ell) = \sum_{\text{state}} A_{n_1}^{\text{tree}}(\ell) \dots A_{n_4}^{\text{tree}}(\ell)$$

$$\bar{C}_{i_1 i_2 i_3}(\ell) = (\dots)_3 - \sum_{j \neq \dots} \frac{J_{i_1 i_2 i_3 j}}{D_j}$$

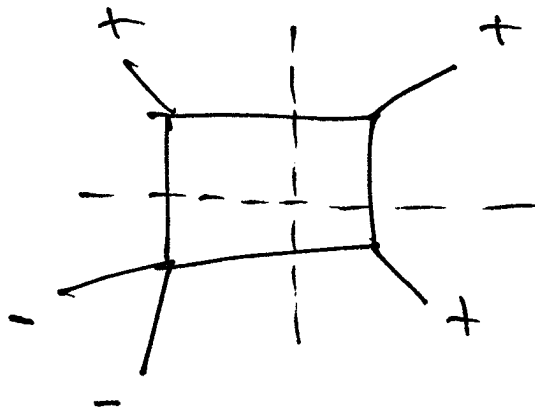
EXAMPLE - concretely

eg. $A_5^{1\text{-loop}}(1^-, 2^-, 3^+, 4^+, 5^+)$

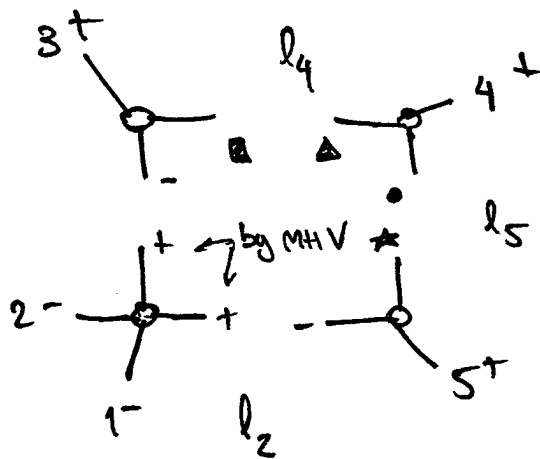
↑ @ tree level we know, Berke-Taylor



getting d_1 :



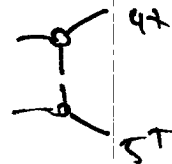
PUT EVERYTHING ON SHELL



- + ±
- ▲ - ±
- ± -
- ★ ± +

some choices
(check these)

only one survives by MHV rules
on one tree graph



once again shift involved, eg

$$d_4 = \frac{z}{2} \langle 3^+ | \sigma^+ | 4^+ \rangle$$

⋮

same trick as tree.

Computing this is the discontinuity
@ integrand level

$$\text{Diagram} = \frac{-\frac{1}{2} \langle 12 \rangle^3 S_{34} S_{45}}{\langle 23 \rangle \langle 34 \rangle \langle 45 \rangle \langle 51 \rangle} \leftarrow \text{eg } (k_4 + k_5)^2$$